

TSC Simulation of Disruptive Current Termination on JT-60U Reversed Shear Plasmas

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Motivation for Study on Disruption Dynamics

- Disruptive termination has been studied and clarified.

e.g. Vertical Displacement Event avoidance at neutral point

(Y. Nakamura et al., Nucl. Fusion 36, 643 (1996).)

(R. Yoshino et al., Nucl. Fusion 36, 295 (1996).)

Measurements of **positive current spike** (I_p spike) with a **flattening of current profile** in DIII-D disruptive Normal Shear (**NrS**) plasmas.

(P.L. Taylor et al., Phys. Rev. Lett. 76, 916 (1996).)



Contrary to common belief of occurring **negative I_p spike** in Reversed Shear (**RS**) plasmas, **positive I_p spike** has been observed in JT-60U.



- Disruptions of **RS** plasmas have not been well understood.

Understanding of disruption dynamics of **RS/NrS** plasmas is important for long-pulsed tokamak reactors like ITER.

Contents

- Various I_p spikes, **positive I_p spike** or **negative I_p spike**, was observed at the thermal quench in JT-60U experiments.

The mechanisms has not been well understood.

- Tokamak Simulation Code (TSC) can model realistic JT-60U disruptions which simulates experiments, incorporating such physics as:
 - Effects of pressure drop (**β_p drop**) including shell structure outside plasmas
 - Dynamics of **rapid I_p profile change** inside plasmas

➔ A new understanding explicates the I_p spike behavior of tokamak disruptions in detail .

JT-60U Experiment in contrast to Analytical Model

- A radial shift of plasma through β_p drop results in a **positive I_p spike** during a thermal quench phase.

$$\delta \tilde{I}_p = - \frac{\delta \tilde{\beta}_p}{2 \left[\ln\left(\frac{8R_p}{a}\right) + \frac{1}{2}(l_i - 1) - 2 \left\{ (n-1) \left(\ln\left(\frac{8R_p}{a}\right) + \frac{1}{2}(l_i - 3) + \beta_p \right) + 1 \right\} \right]}$$

$$\approx -0.05 \delta \tilde{\beta}_p \quad (\delta \tilde{\beta}_p < 0) \quad (\because \delta \psi = 0, \delta F_R = 0) \quad n: \text{decay index}$$

(A. Fukuyama et al., Jpn. J. Appl. Phys. 14, 871 (1975).)



$\delta \tilde{I}_p$ depends very weakly on R_p ($\sim R_J$).

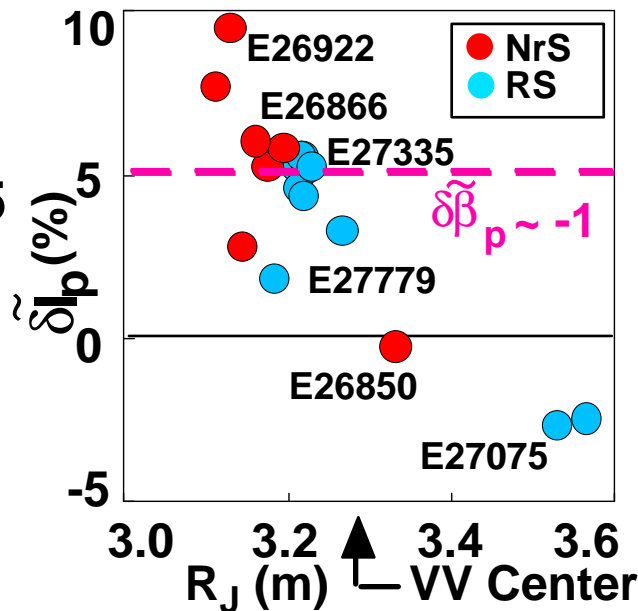
Surprising feature of I_p spike

$I_p \sim 2-3\text{MA}$

$\beta_p \sim 1$

$q_{\text{eff}} \sim 3.5-4.5$
(**NrS**)

$q_{\text{eff}} \sim 4.5-7$
(**RS**)



- $\delta \tilde{I}_p (= \delta I_p / I_p)$ with inner I_p center (R_J) was larger than that of outer R_J .

- Maximum $\delta \tilde{I}_p$ ($\sim 10\%$) was observed in **NrS** discharges at smaller R_J .

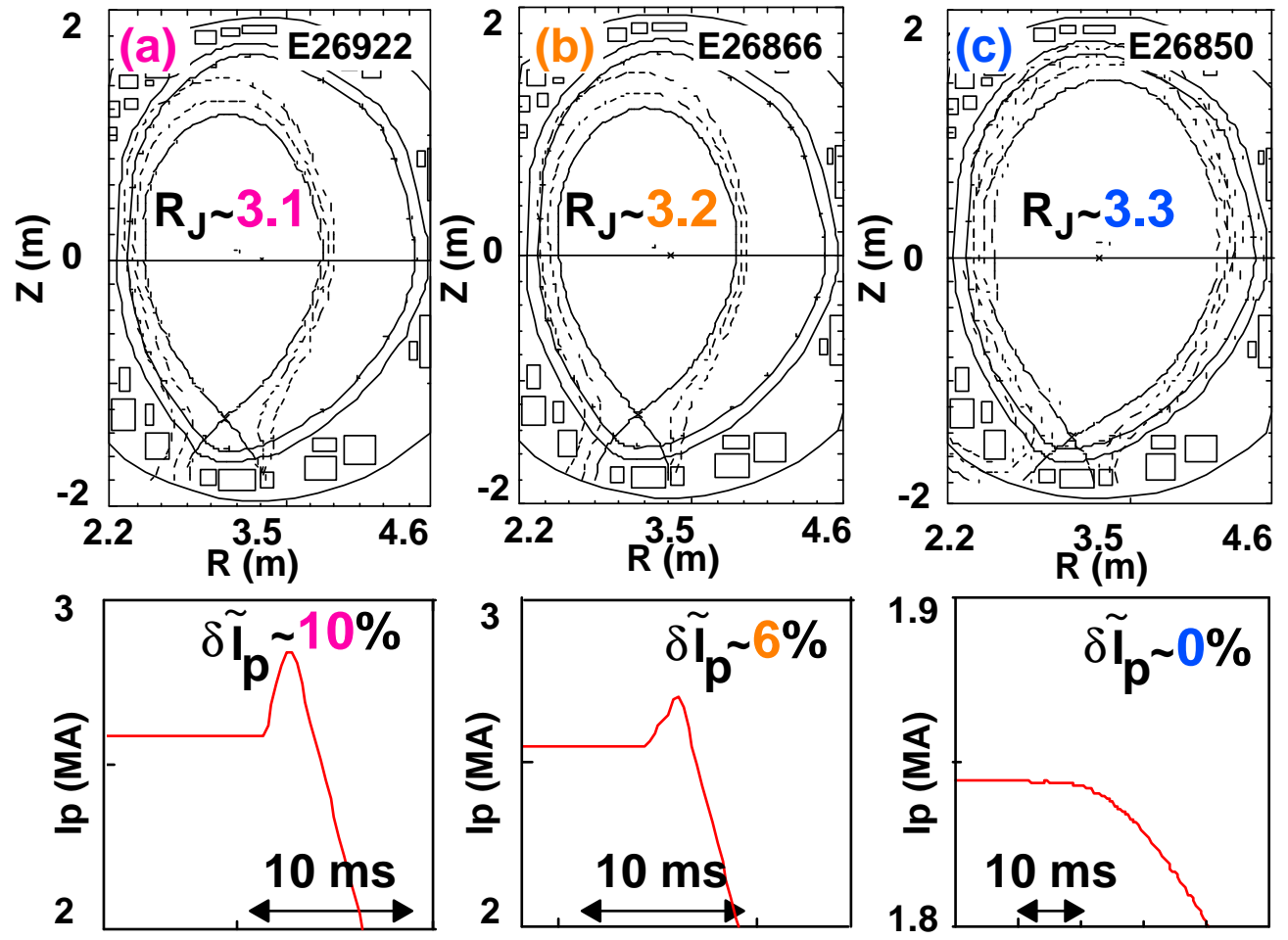
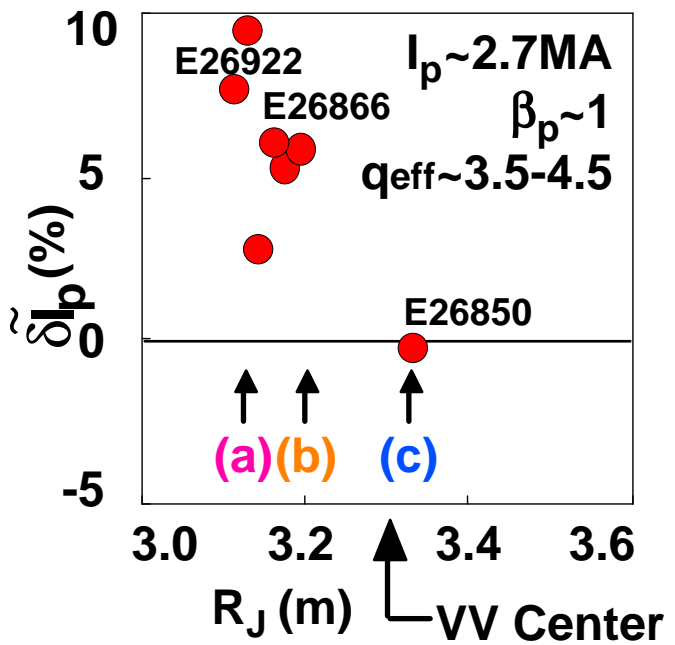
- Negative I_p spike ($\sim -5\%$) was observed in **RS** discharges at larger R_J .



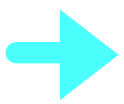
It is found that $\delta \tilde{I}_p$ depends on R_J

Current Spike Observed during High β_p NrS Disruption

$\delta \tilde{I}_p$ depends on R_J .



$\delta \tilde{I}_p$ with **smaller** R_J
 $>$ $\delta \tilde{I}_p$ with **larger** R_J



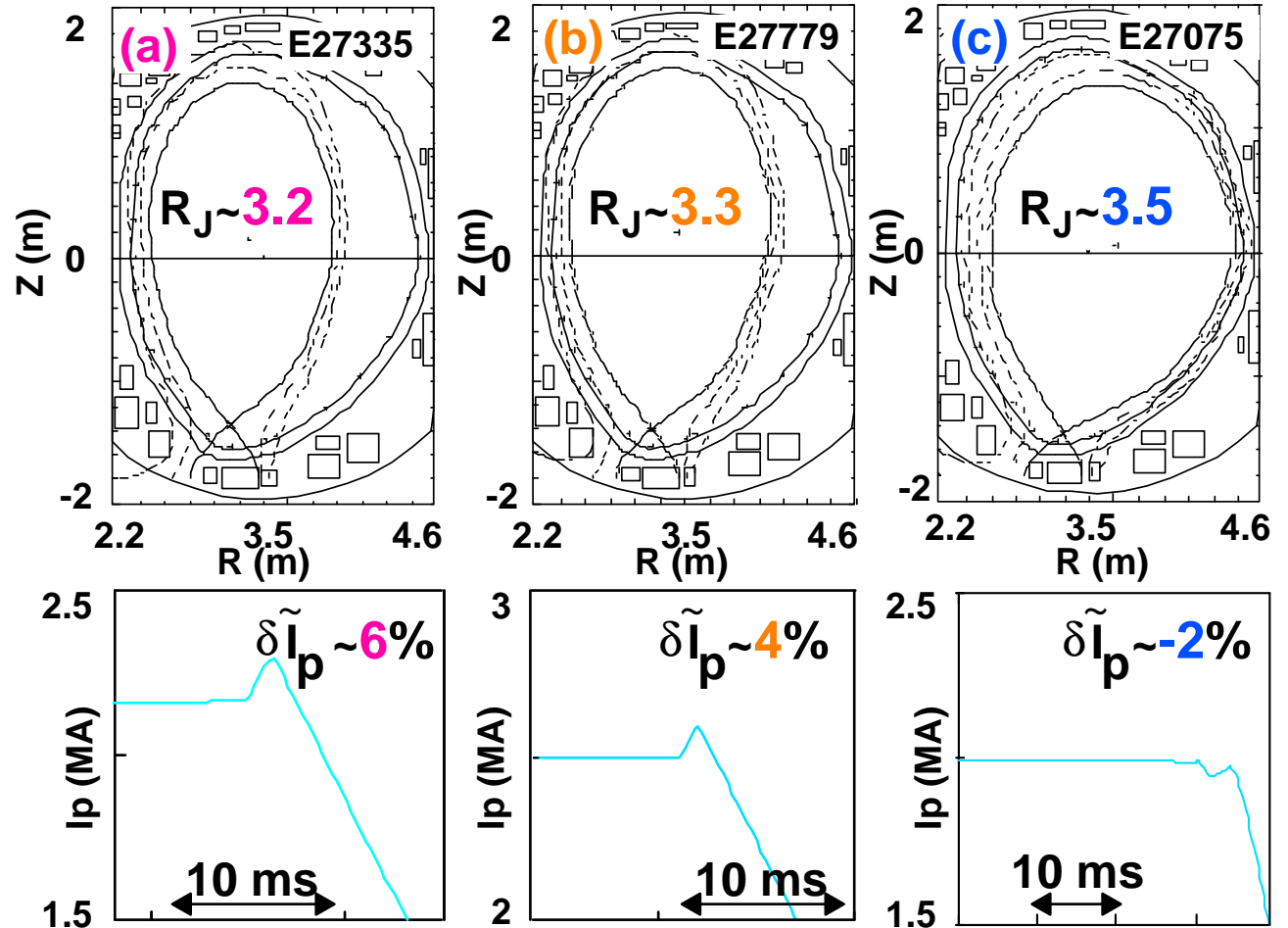
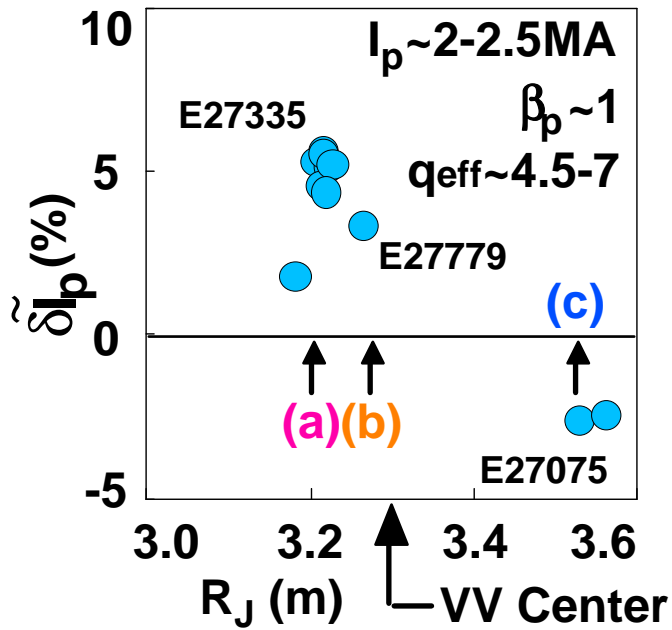
$\delta \tilde{I}_p \sim 10\%$
 with $R_J \sim 3.1$ m

$\delta \tilde{I}_p \sim 6\%$
 with $R_J \sim 3.2$ m

$\delta \tilde{I}_p \sim 0\%$
 with $R_J \sim 3.3$ m

Current Spike Observed during High β_p RS Disruption

$\delta \tilde{I}_p$ depends on R_J .



$\delta \tilde{I}_p$ with **smaller** R_J
 $>$ $\delta \tilde{I}_p$ with **larger** R_J

$\delta \tilde{I}_p \sim 6\%$
 with $R_J \sim 3.2 \text{ m}$

$\delta \tilde{I}_p \sim 4\%$
 with $R_J \sim 3.3 \text{ m}$

$\delta \tilde{I}_p \sim -2\%$
 with $R_J \sim 3.5 \text{ m}$

Shell effect on the current spike under the β_p drop ? - - - \rightarrow TSC simulation

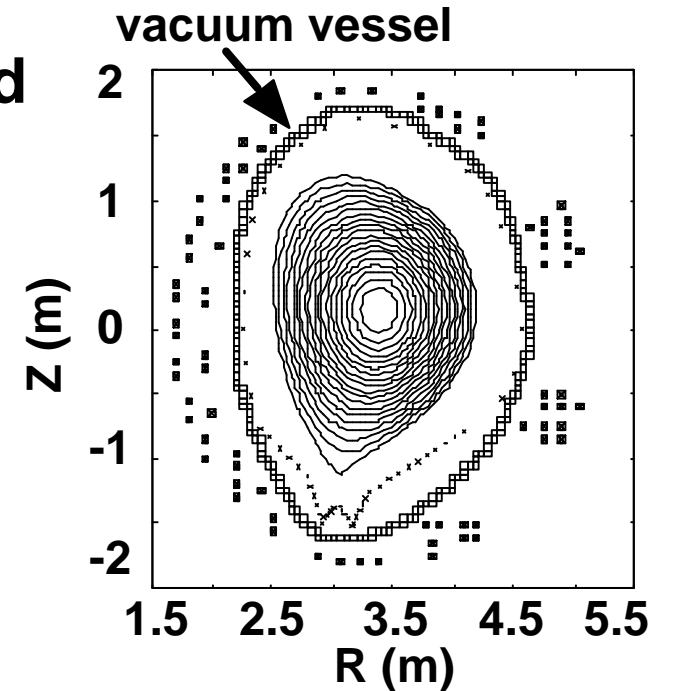
TSC Modeling of Axisymmetric MHD

- Usual set of MHD eqs. with modified momentum eq.

$$\frac{\partial \mathbf{m}}{\partial t} + \mathbf{F}_V(\mathbf{m}) = \mathbf{j} \times \mathbf{B} - \nabla p$$

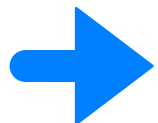
\mathbf{m} : plasma momentum density ($= \rho \mathbf{v}$)

$\mathbf{F}_V(\mathbf{m})$: artificial viscosity operator
instead of $\mathbf{v} \cdot \nabla \mathbf{v}$



TSC configuration of JT-60U

- Free-boundary axisymmetric MHD → Realistic plasma configuration (TSC, DINA)
- Disruption dynamics including shell structure → β_p drop (TSC, DINA, ASTRA)
- Resistive diffusion of I_p profile → Current flattening (TSC etc.)



TSC enables us to realistically simulate with experiments.

Current Spike due to β_p Drop at Outer Position ($R_J=3.8m$)

TSC

β_p drop ($\delta\tilde{\beta}_p \sim -1$) leads to inward shift ($\delta R_J \sim -0.2m$) and results in $\delta\tilde{I}_p \sim -3\%$.

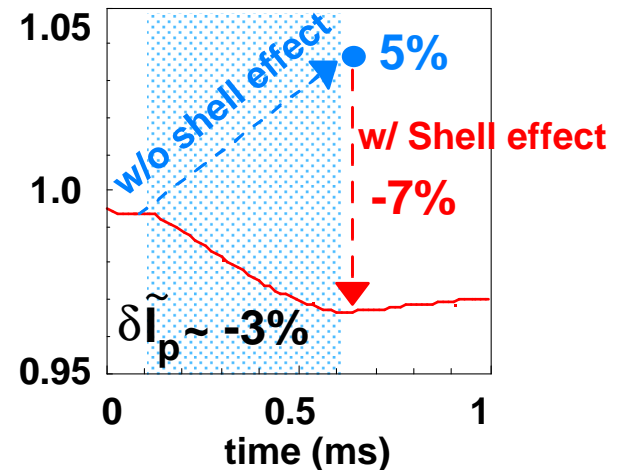
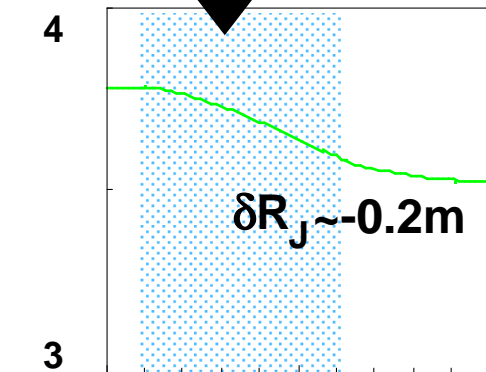
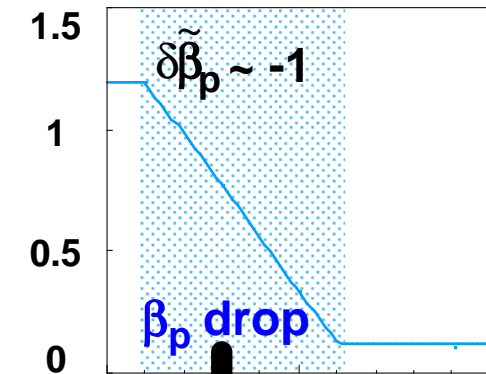
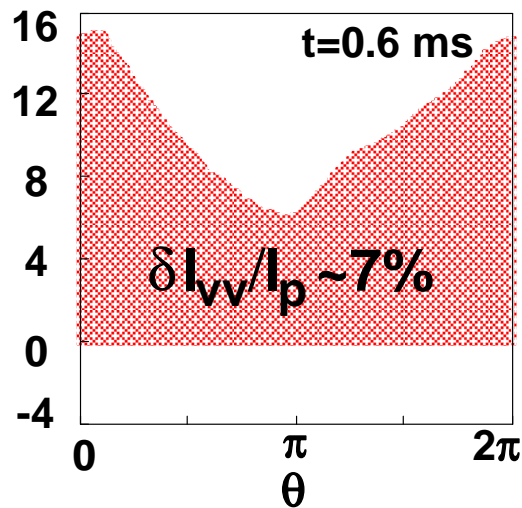
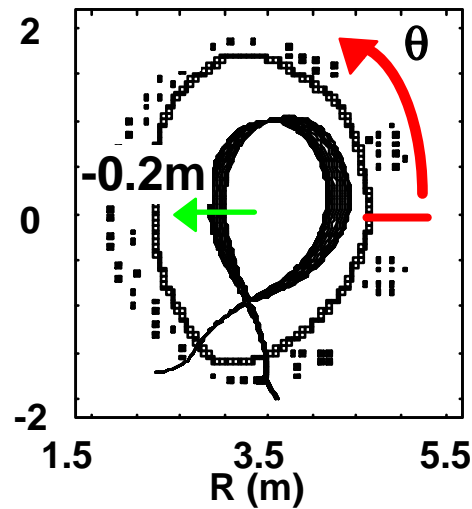
- Analytical model without shell effects

The β_p drop results in a positive I_p spike.

$$\delta\tilde{I}_p \approx -0.05 \delta\tilde{\beta}_p \sim 5\%$$

- TSC simulation

Induced **toroidal eddy current** in the vacuum vessel



Negative I_p spike

β_p drop ($\delta\tilde{\beta}_p = -0.5$) leads to inward shift ($\delta R_J \sim -0.1\text{ m}$) and results in $\delta\tilde{I}_p \sim 1\%$.

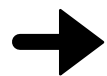
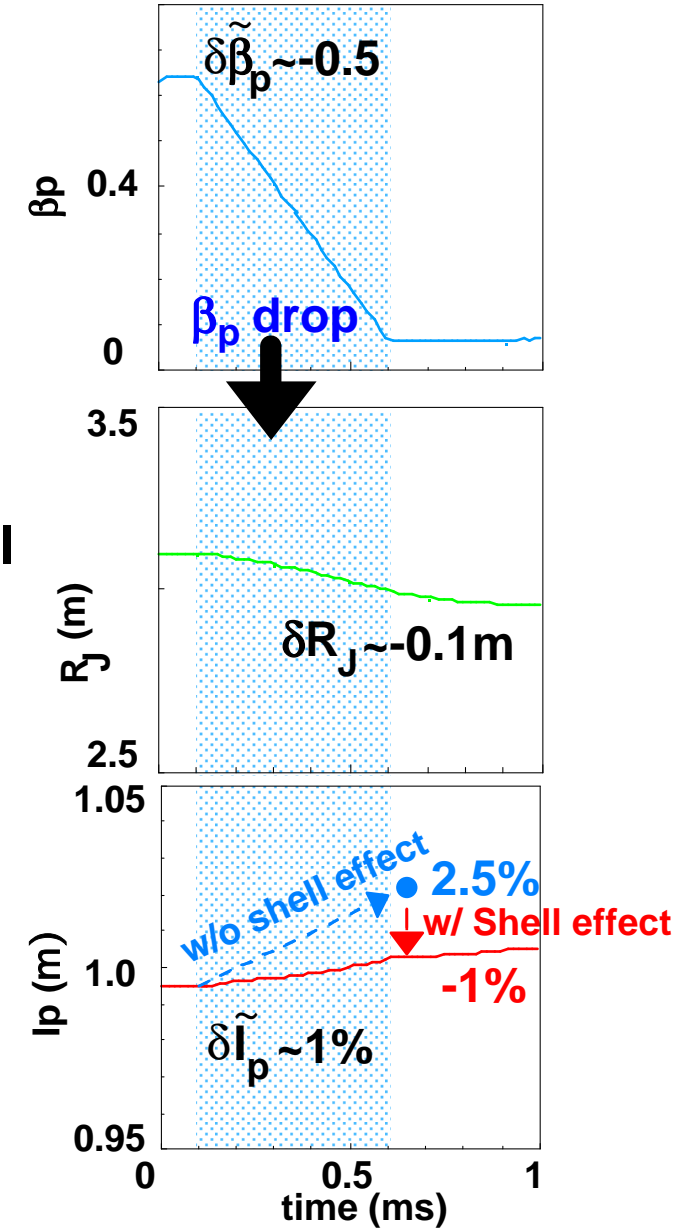
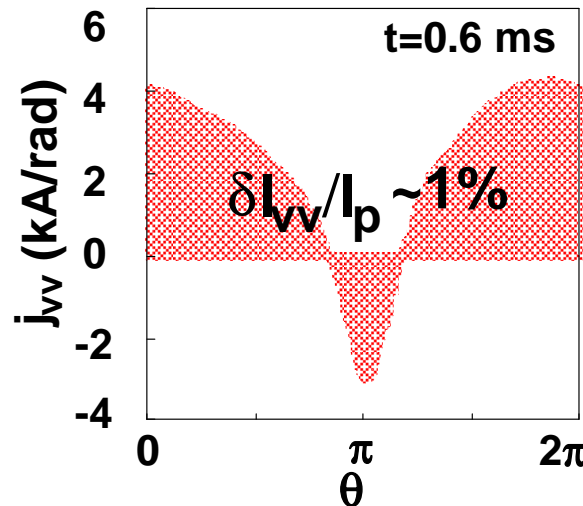
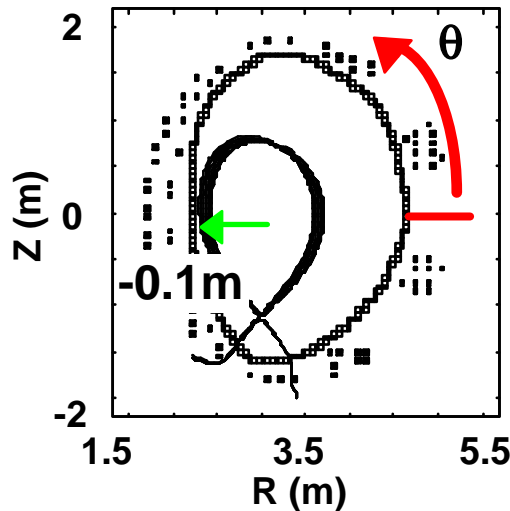
- Analytical model without shell effects

The β_p drop results in a positive current spike.

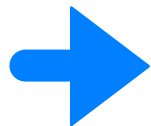
$$\delta\tilde{I}_p \approx -0.05 \delta\tilde{\beta}_p \sim 2.5\%$$

- TSC simulation

Induced **toroidal eddy current** in the vacuum vessel



Increasing I_p



The effect of β_p drop results in the R_J dependence of $\delta\tilde{I}_p$.

Abrupt Change in Current Profile

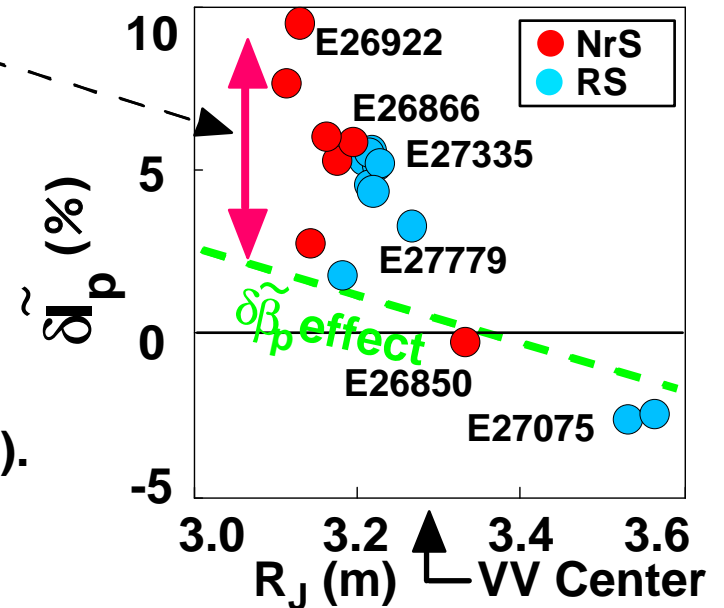
$\delta\tilde{\beta}_p$ effect is insufficient to explain $\delta\tilde{I}_p$.

Another mechanism?

Measurements of **current flattening** in DIII-D **NrS** plasmas

P.L. Taylor et al., Phys. Rev. Lett. 76, 916 (1996).

Effect of **abrupt change in current profile**



Abrupt current profile change model in TSC

- Ohm's Law modified by **Hyper-Resistivity**

(A.H. Boozer, J. Plasma Phys. 35, 133 (1986).)

- Destruction of magnetic surface (tearing mode)
- Dissipating energy and conserving helicity

$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \eta \mathbf{j} - \frac{\mathbf{B}}{B^2} \nabla \cdot \left(\lambda \frac{\mathbf{j}_{\parallel}}{B} \right) \quad \lambda: \text{arbitrary positive function}$$

→ TSC models anomalous current viscosity

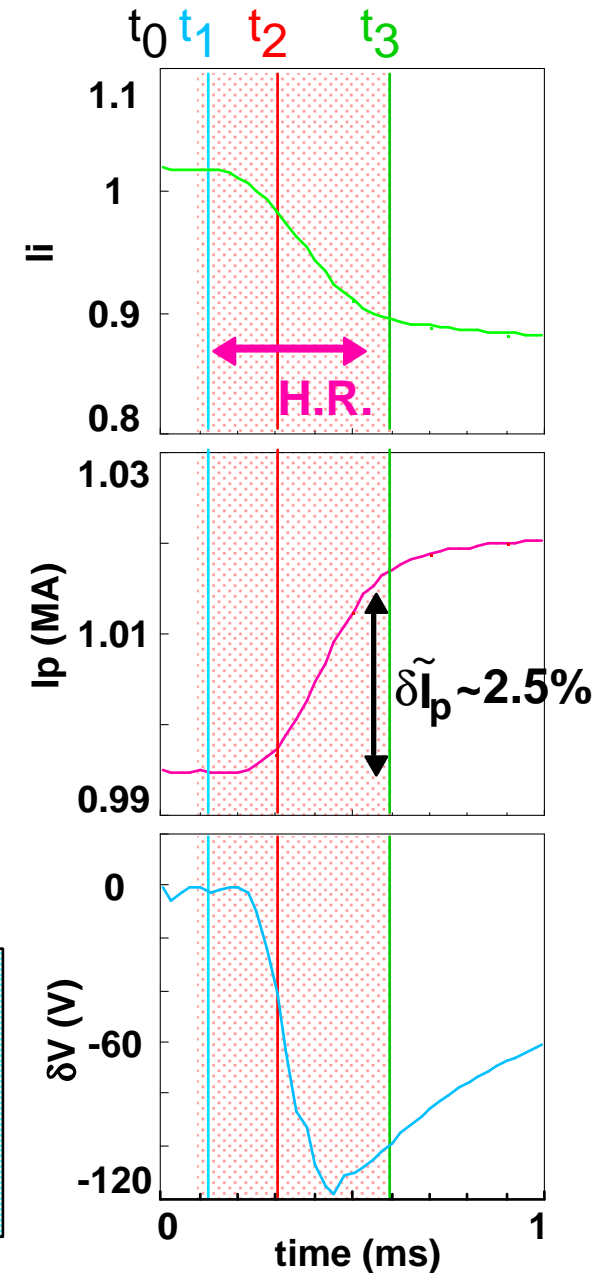
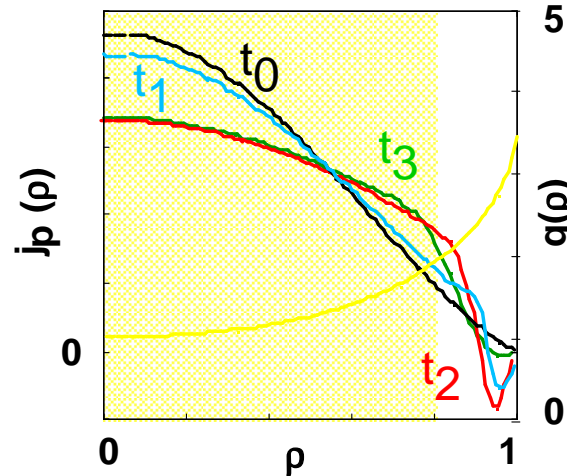
D.J. Ward, S.C. Jardin, Nucl. Fusion 29, 905 (1996).

Current Profile Change in **NrS** plasmas ($\beta_p \sim 0$) **TSC**

TSC with the hyper-resistivity reproduced a redistribution in current profile

- Hyper-Resistivity causes a **lowering of I_i** .

→ **Positive I_p spike**
 $\delta \tilde{I}_p \sim 2.5\%$
→ **Negative voltage spike**
 $\delta V \sim -120V$



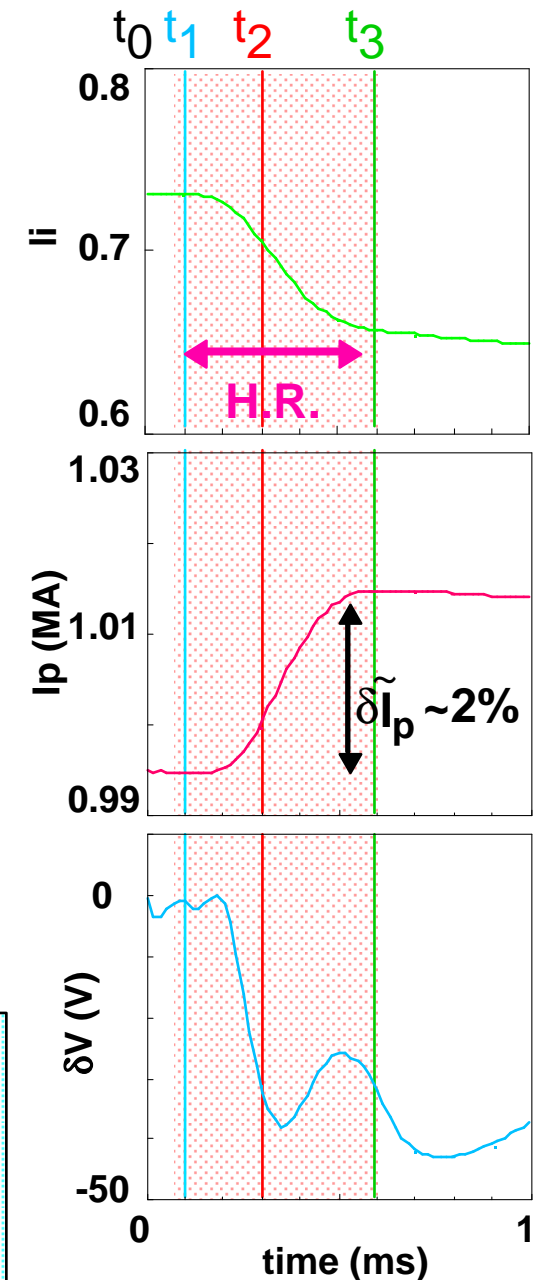
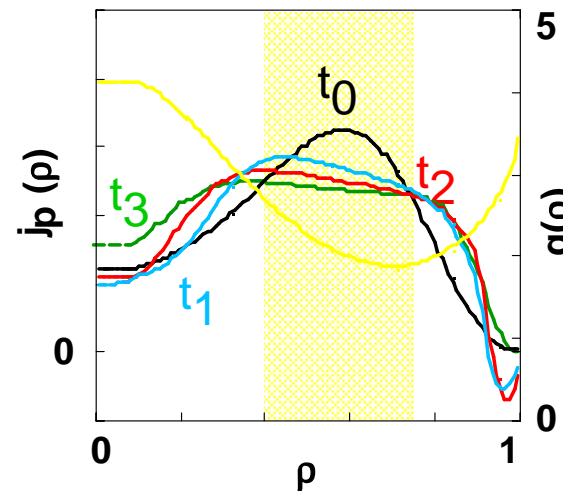
➔ Current flattening explicates both spikes of **positive I_p** and **negative voltage** similar to experiments.

Current Profile Change in RS plasmas ($\beta_p \sim 0$) TSC

Redistribution of current profile in the same manner as **NrS** plasmas

- Hyper-Resistivity causes a **further lowering of I_i** .

▶ **Positive I_p spike**
 $\delta \tilde{I}_p \sim 2\%$
▶ **Negative voltage spike**
 $\delta V \sim -40V$



▶ **Further lowering of I_i** results in both spikes
of **positive I_p** and **negative voltage** even in
RS plasmas.

Conclusions

- Important effects on I_p spikes

- (a) Induced eddy current in vacuum vessel

β_p drop causes the I_p spike that is positive with smaller R_J while negative with larger R_J . (strong R_J dependence)

- (b) Abrupt current profile change

Further lowering of I_i similar to NrS plasmas is a plausible candidate for the mechanism of positive I_p spike observed even in RS plasmas.

- A new understanding obtained by TSC simulation reasons out the various current spike behaviors of JT-60U RS/NrS plasmas.

