TRANSPORT AND STABILITY IMPLICATIONS FOR SHAPE AND ASPECT RATIO OF STEADY-STATE, HIGH-PERFORMANCE TOKAMAKS

by C.C. Petty General Atomics

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OUTLINE

- I. Optimum tokamak study by Lin-Liu/Stambaugh
- II. Comparison with similar study by Menard
- III. Transport dependence on shape (A, κ)



- Lin-Liu/Stambaugh constructed equilibria with
 - Bootstrap fraction of 99%, fully aligned
 - P' = 0 at separatrix
 - Broad, nearly optimal, pressure profile
 - ★ Edge ITB?
- Ideal ballooning β limit found using BALOO
 - Bulk of plasma has second stability access
 - Ballooning limit occurs at a point near edge
 - Wall stabilization assumed for kinks
- Systematic shape study spanned
 - 1.5 $\leq \kappa \leq$ 6.0
 - $-1.2 \le A \le 7.0$



HIGH BETA, HIGH ELONGATION, HIGH BOOTSTRAP EQUILIBRIUM

• A = 1.6, κ = 4.0, δ = 0.5, β_{T} = 73%, β_{N} = 8.0, β_{P} = 1.6





 β_N is optimal at $\kappa = 3-4$

 β_N dependence close to A^{-1/2}





STRONG SHAPING (δ) IS NEEDED TO TAKE FULL ADVANTAGE OF HIGH ELONGATION

• Beta increases with δ for $\kappa \ge 3$





• Trade-off between fusion power and bootstrap current at a given normalized beta

$$\beta_{T} \beta_{P} = 25 \left(\frac{1+\kappa^{2}}{2}\right) \left(\frac{\beta_{N}}{100}\right)^{2}$$

$$f_{BS} = C_{BS} \beta_{P} / \sqrt{A}$$

$$P_{F} \propto \beta_{T}^{2} B^{4}$$

• Shape dependence of ideal ballooning stable beta

$$\beta_{\mathsf{N}} = 10 \ (\mathsf{b}_0 + \mathsf{b}_1 \kappa + \mathsf{b}_2 \kappa^2 + \mathsf{b}_3 \kappa^3) \ \mathsf{coth} \left(\frac{\mathsf{d}_0 + \mathsf{d}_1 \kappa}{\mathsf{A}^{\mathsf{m}}}\right) \frac{1}{\mathsf{A}^{\mathsf{n}}}$$

b ₀	-0.7748	d ₀	1.8524
b ₁	1.2869	d ₁	0.2319
b ₂	-0.2921	m	0.6163
b ₃	0.0197	n	0.5523



II. IDEAL WITH-WALL BALLOONING LIMIT FOR FULLY SELF-SUSTAINED EQUILIBRIA NEARLY SAME BETWEEN MENARD'S AND LIN-LIU'S STUDIES





• While empirical confinement scaling relations of the form

IPB98(y,2) τ =0.0562 II0.93n0.15P<-0.69</th>R1.97m0.19 $\kappa^{0.78}$ A-0.58EGB τ =0.028 I0.83n0.49B0.07P<-0.55</td>R2.11m0.14 $\kappa^{0.75}$ A-0.3

are fully predictive, the trade-offs between different parameters due to operational constraints are not readily apparent

 $\bigstar\,$ e.g., safety factor constraint relates I, $\kappa,$ A, and R

- Casting confinement scaling relations in terms of dimensionless parameters allows the shape and aspect ratio dependences to be easily determined once the operational constraints are specified
 - Can choose kinetic plasma physics parameters like ρ_{\star} and ν_{\star}
 - Can choose MHD parameters like q and β_{N}



DEPENDENCE OF TRANSPORT ON q AND κ

- Experiments on DIII–D resolved the ambiguity between the κ and q scalings of transport by comparing
 - q scan at fixed κ
 - κ scan at fixed q
 - κ scan at fixed I
- For H–mode plasmas, the change in confinement for the above three scans was explained by the unified scaling

$$B\tau \propto q_{95}^{-1.4\pm0.6} \quad \stackrel{2.2\pm0.6}{\kappa}$$

- Note that the q and κ scalings of normalized confinement are different than the I_p and κ scalings of τ
 - ★ Converting dimensionless parameter scalings for H-mode plasmas on DIII-D to engineering parameter scalings gives

$$\tau \propto \mbox{I}_p^{0.76 \pm 0.14} \ \ \kappa^{0.65 \pm 0.16}$$



MEASURED q AND κ SCALINGS OF H–MODE TRANSPORT ARE WEAKER THAN PREDICTION FROM IPB98(y,2)





OPERATIONAL CONSTRAINTS ARE CRITICAL WHEN PROJECTING ASPECT RATIO SCALING OF TRANSPORT

- Aspect ratio affects many important dimensionless parameters — q, β_N , ν_* , f_{BS}, etc.
- Future experiments between DIII–D and NSTX/MAST will directly measure aspect ratio scaling of transport
- For steady-state, high-performance tokamaks, the aspect ratio scaling of confinement is more easily projected by substituting operational constraints for engineering parameters in scaling relations





INCLUDING EFFECT OF ASPECT RATIO ON β_{N}, κ CHANGES R/a DEPENDENCE OF NORMALIZED CONFINEMENT

Confinement scaling relations converted to dimensionless parameters

IPB98(y,2) $B\tau \propto \rho_{\star}^{-2.7} A^{1.2} \kappa^{2.4} \beta_{N}^{1.2} f_{BS}^{-2.1} f_{GR}^{0.0}$

 $\label{eq:EGB} \textbf{EGB} \qquad \textbf{B} \tau \propto \rho_{\star}^{-3.3} \, \textbf{A}^{1.8} \; \kappa^{2.2} \; \beta_{N}^{2.1} \; \textbf{f}_{BS}^{-1.7} \; \textbf{f}_{GR}^{-0.6}$

• Include optimum tokamak scalings $\beta_N \propto A^{-1/2}$, $\kappa \propto A^{-1/2}$



FUSION GAIN OPTIMIZES FOR A = 2.2–3.0 FOR STEADY-STATE, HIGH-PERFORMANCE TOKAMAKS AT STABILITY LIMIT



• Assume $B = B_c (1-A^{-1})$ where B_c (= field at centerpost) is fixed



RECIRCULATING POWER FRACTION OPTIMIZES FOR A=2.4–2.8 AT STABILITY LIMIT





KEY RESULTS OF TRANSPORT DEPENDENCE ON SHAPE (A, κ)

- Transport dependence on elongation and safety factor are weaker than IPB98(y,2) relation but close to EGB relation
- Transport dependence on aspect ratio is more apparent when the operational constraints (β_N, κ, f_{BS}, f_{GR}) are directly incorporated into the confinement scaling relation
- For "optimum tokamak", fusion gain is optimized between aspect ratio of 2.2 and 3.0 (depending upon which confinement scaling relation is used)
- If stability limit and elongation are assumed independent of aspect ratio, then fusion gain optimizes at higher R/a

