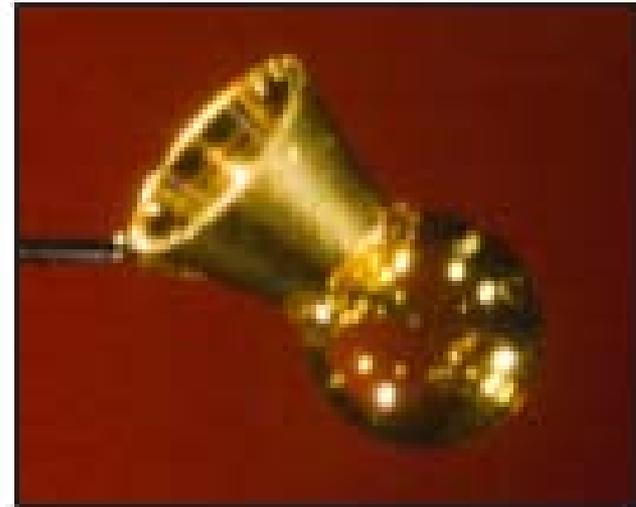
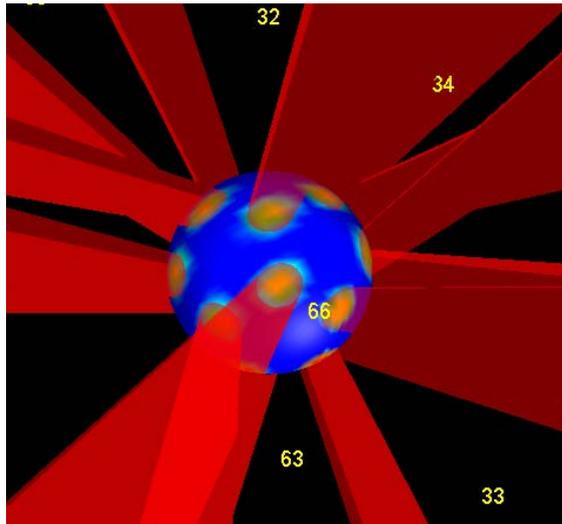


# Advanced ignition options for laser ICF

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R. Betti

*University of Rochester and Princeton Plasma Physics Laboratory  
FPA Meeting, Washington DC, December 1-3, 2010*



# The NIF can explore advanced ignition options



- **With day-one hardware, the NIF can explore high-gain shock ignition**
  - **Polar Shock Ignition (uses half the NIF beams to drive the implosion and the other half to drive the ignitor shock)**
- **Fast Ignition requires major hardware upgrades: 100kJ-class multi-PW laser [also talk by P. Patel at this meeting]**
- **Polar Direct Drive requires minor upgrades: multi-FM or 2D-SSD (talk by J. Soures at this meeting)**

# FSC Collaborators

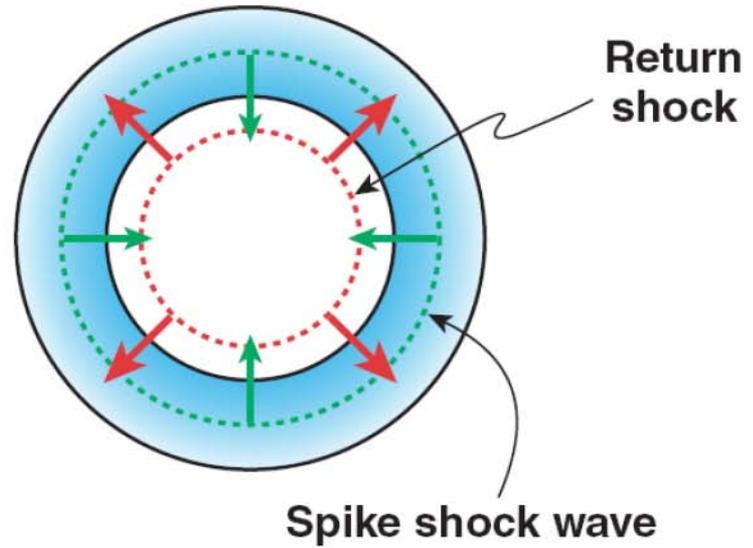
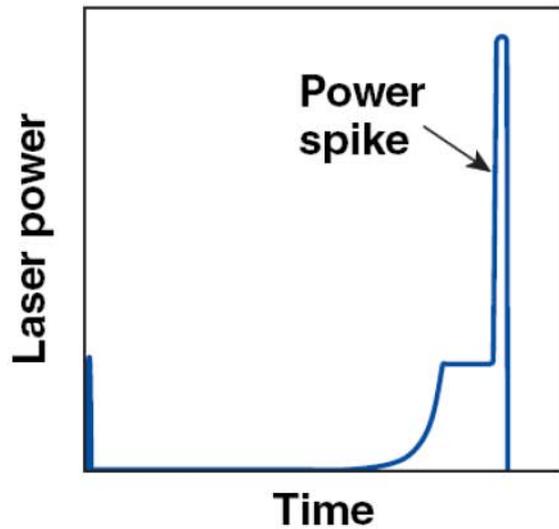


K. Anderson (LLE)  
W. Theobald (LLE)  
R. Nora (LLE)  
C. Stoeckl (LLE)  
M. Hohenberger (LLE)  
A. Solodov (LLE)  
C. Ren (LLE)  
G. Fiksel (LLE)  
PY Chang (LLE)

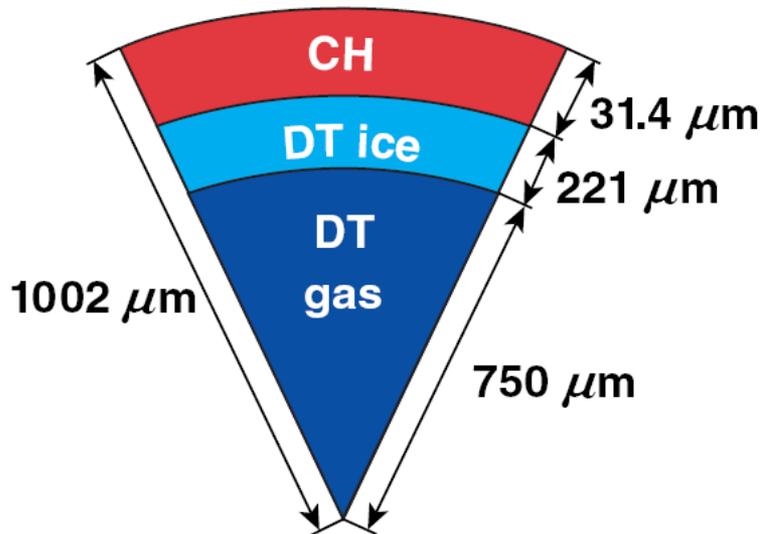
F. Beg (UCSD)  
M. Wei (UCSD)  
R. Stephens (GA)  
P. Patel (LLNL)  
H. McLean (LLNL)  
R. Freeman (OSU)  
L. Van Woerkon (OSU)  
Y. Sentoku (UNR)  
W. Mori (UCLA)

J. Tonge (UCLA)  
R. Petrasso (MIT)  
CK Li (MIT)  
F. Seguin (MIT)  
J. Frenje (MIT)

# Shock Ignition



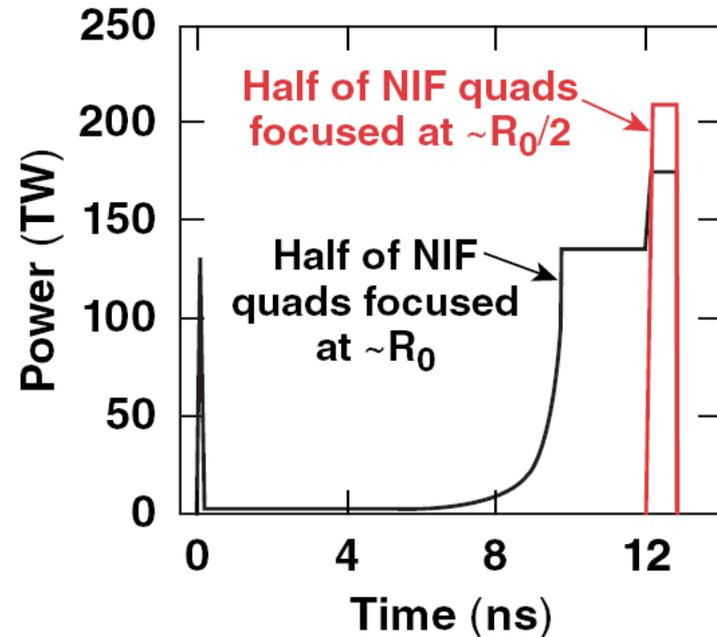
# A thick plastic-ablator shock-ignition target for the NIF has been designed using existing NIF phase plates



Gain (1-D)	70
$\rho R$ (g/cm <sup>2</sup> )	2.6
$V_{\text{imp}}$ (μm/ns)	300
IFAR <sub>2/3</sub>	30

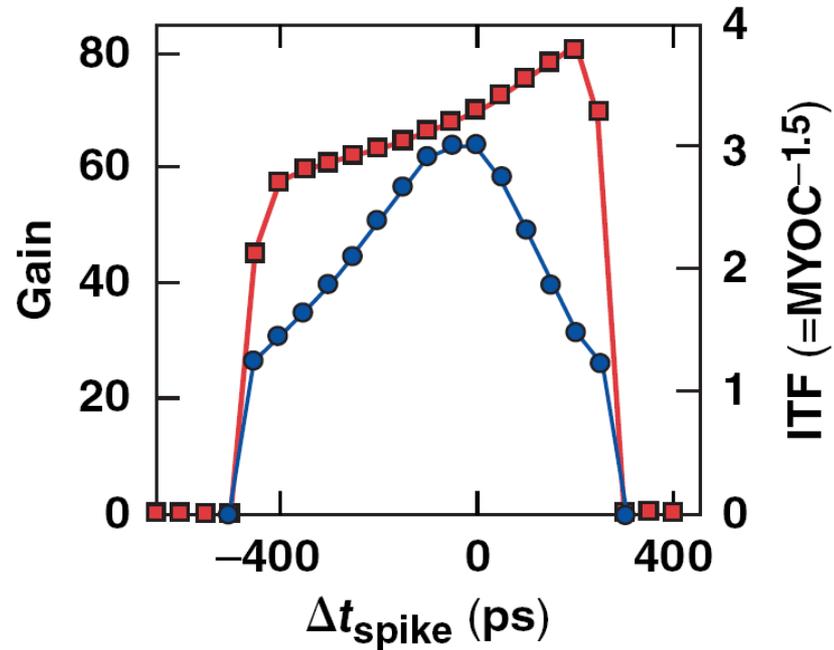
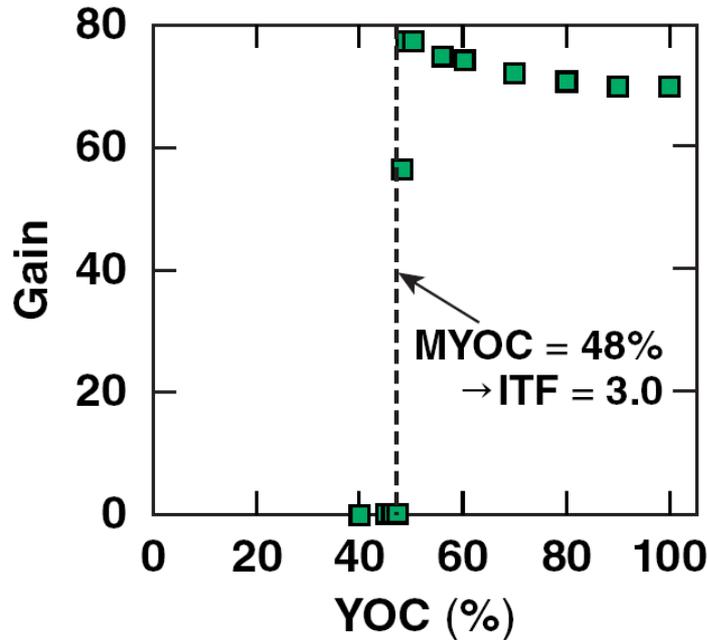
$$\text{IFAR}_{2/3} = \frac{R}{\Delta R} \text{ at } R = \frac{2}{3}R_0$$

Input energy: 680 kJ



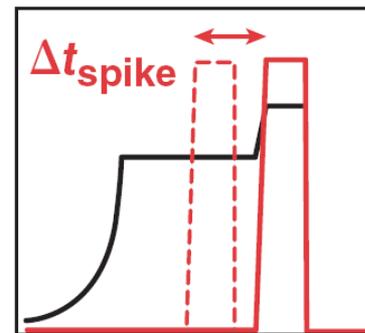
1-D beam profiles approximate polar drive.\*

# Plastic-ablator shock-ignition targets are robust to shock timing and reduced clean volumes

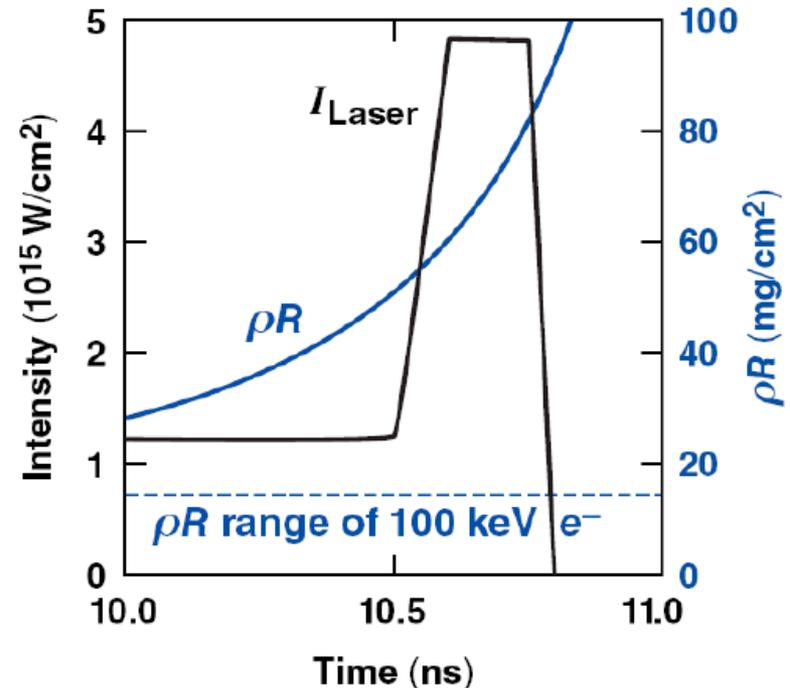
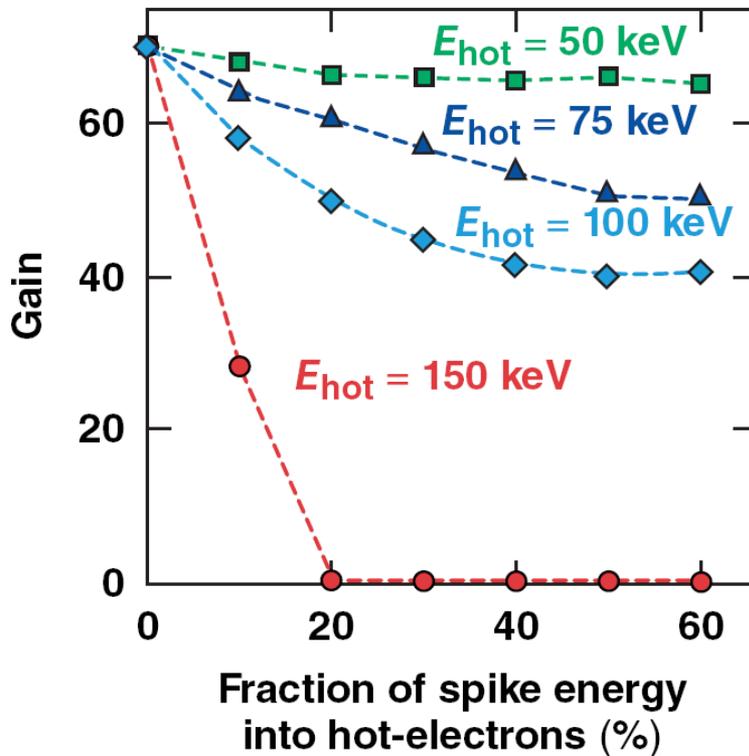


**ITF for indirect-drive point design\* is ~5.3 (MYOC = 33%) at 1 MJ.**

\*J. Lindl, presented to the JASON Review Committee Study #JSR-09-330, San Diego, CA, 14–16 January 2009.



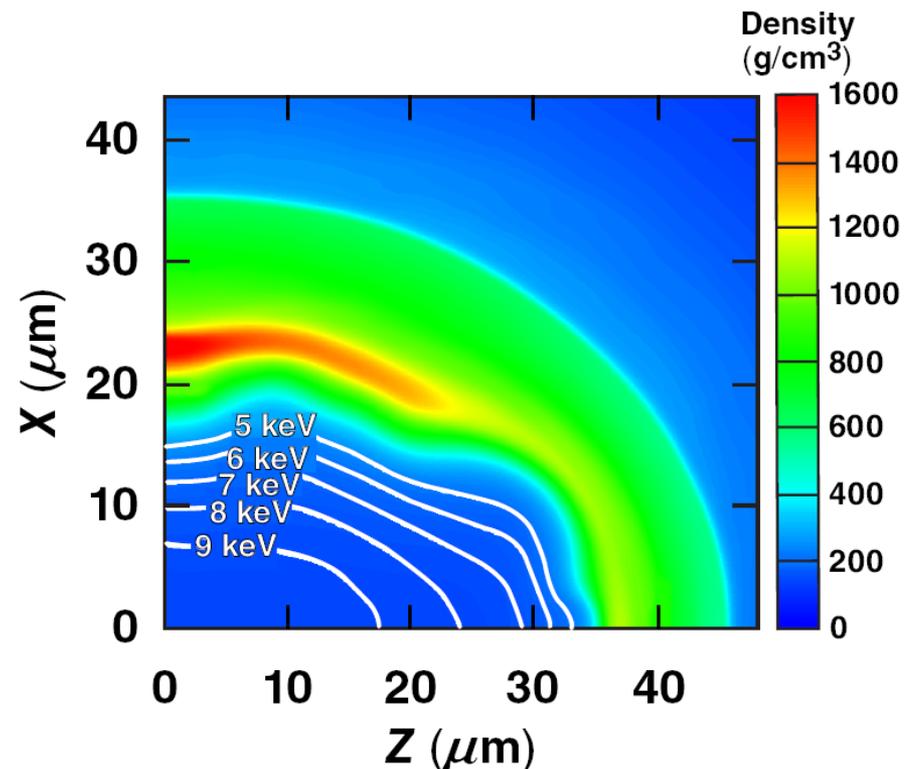
# The plastic-ablator SI design is robust to hot electrons up to 100 keV at 60% of laser energy during the spike pulse



# Symmetric 2-D *DRACO* simulations performed with similar targets indicate robustness to ice roughness $>3.5\text{-}\mu\text{m rms}$



- Symmetric laser irradiation
- *DRACO* simulations with  $3.5\text{-}\mu\text{m-rms}$  roughness in modes  $\ell = 2$  to 50
- Target ignites with full gain
- Upper limit on robustness to ice modes not yet explored
- Other nonuniformity studies to follow (imprint, target offset, polar drive, etc.)

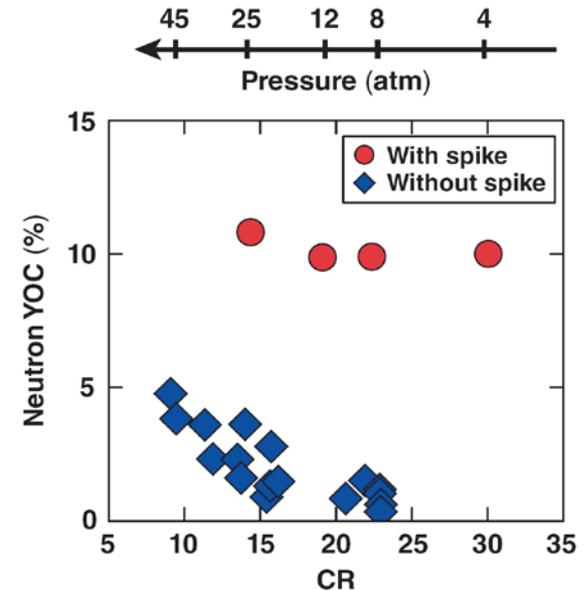
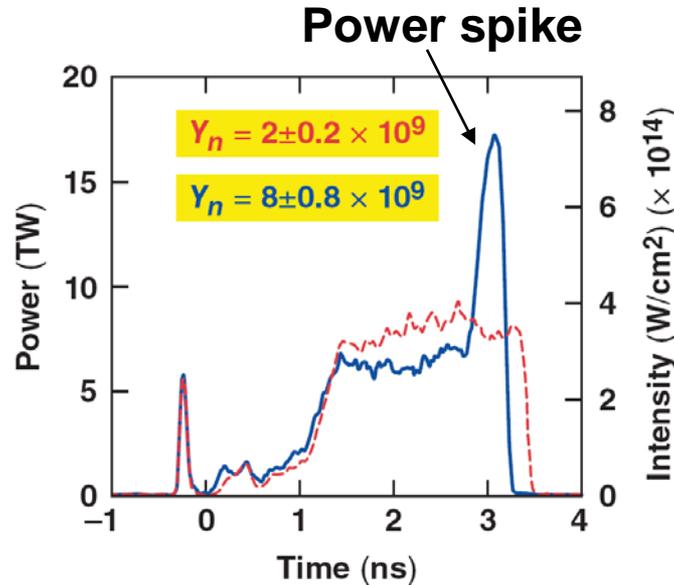
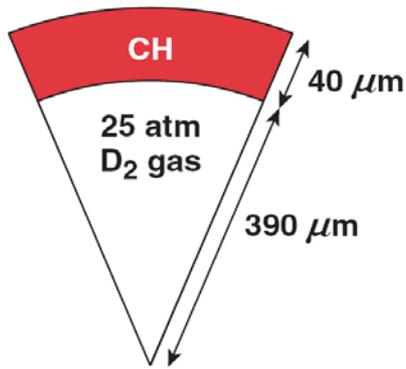


**A Schmitt (NRL) has also developed robust sub-MJ high-gain designs for KrF**

# Shock-ignition experiments on OMEGA have shown improved performance when a shock launching power spike is added at the end of the laser pulse



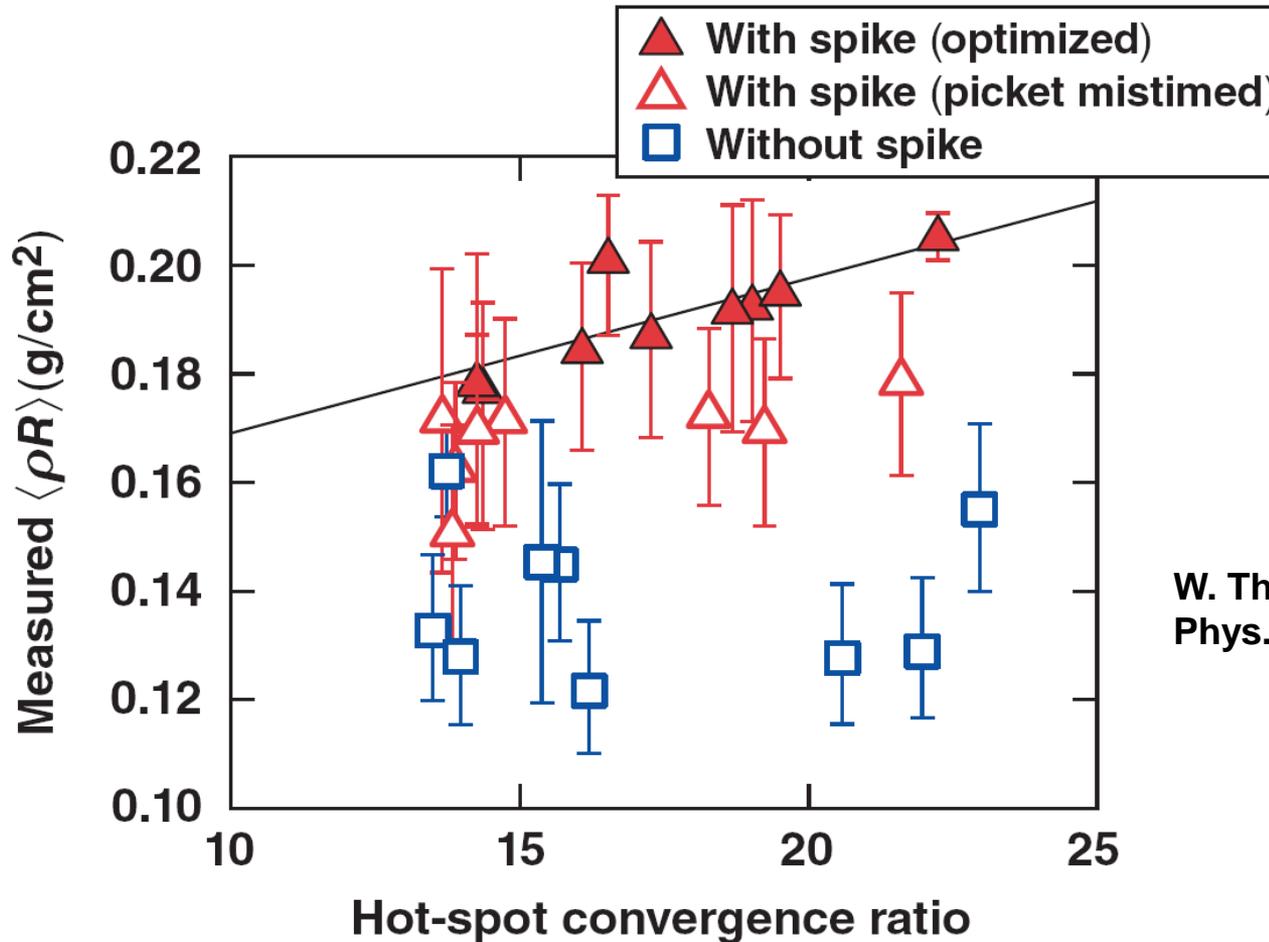
$E_L = 19 \text{ kJ}$ ,  $\alpha = 1.3$ ,  
 $V_j = 1.7 \times 10^7 \text{ cm/s}$ , SSD off



The neutron yield increases considerably when a shock is launched at the end of the pulse.

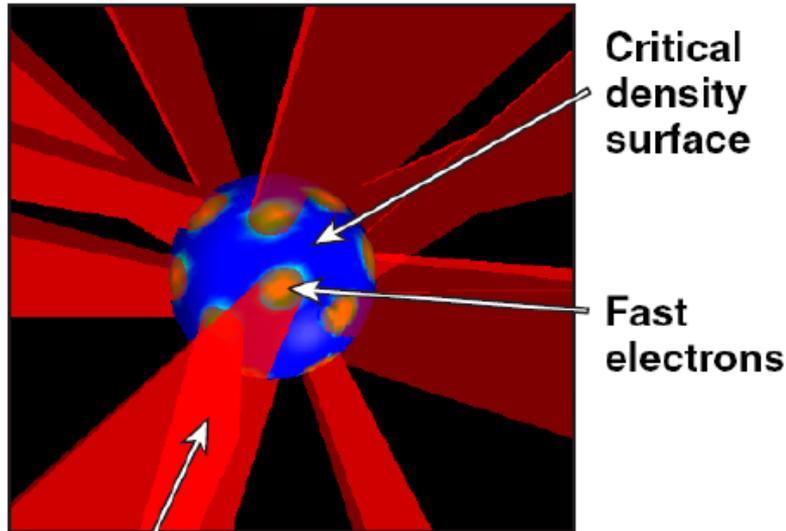
The measured-to-calculated neutron-yield ratios are close to 10% for a hot-spot convergence ratio of 30.

# Higher $\langle \rho R \rangle$ exceeding = 0.2 g/cm<sup>2</sup> where measured in implosions with late spike

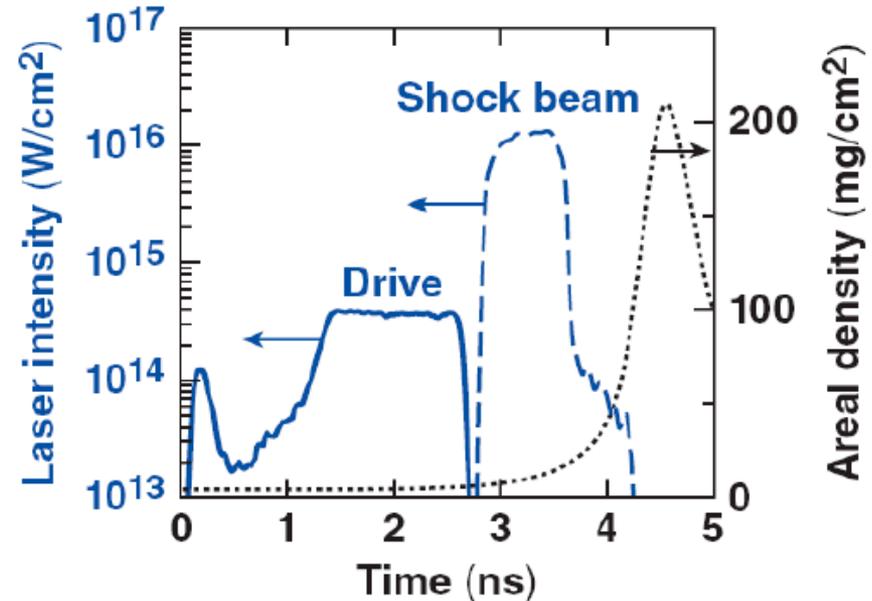


The shock-ignition pulse-shape implosions show an improved performance with respect to compression and neutron yields.

# 60 OMEGA beams were split into 40 low-intensity drive beams and 20 tightly focused beams to study LPI in shock ignition

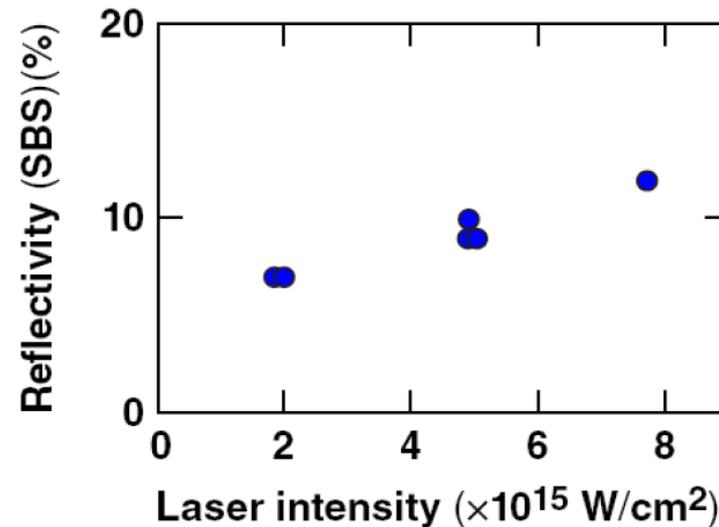
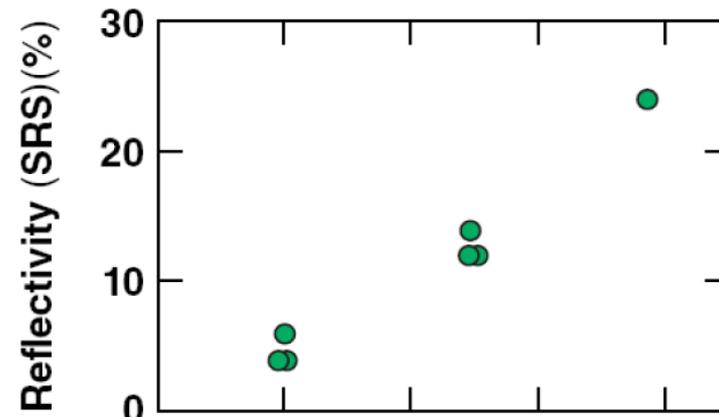
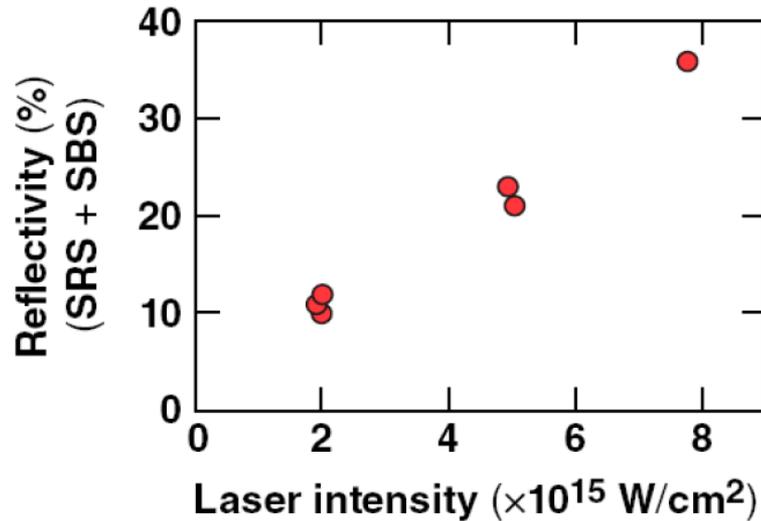


Shock beams  
 $\sim 10^{16}$  W/cm<sup>2</sup>



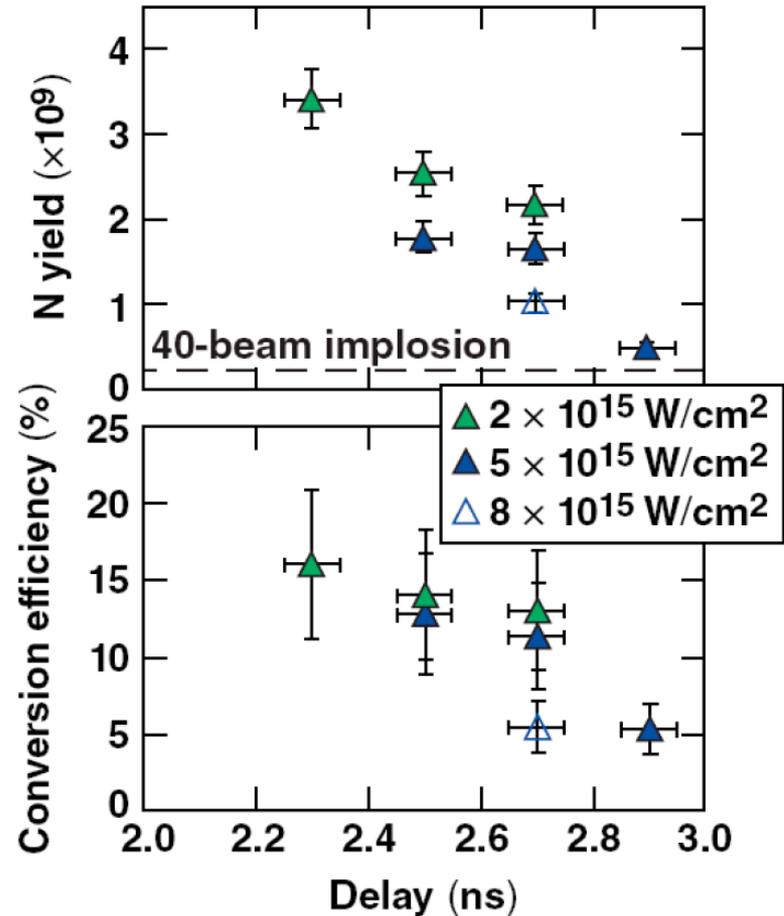
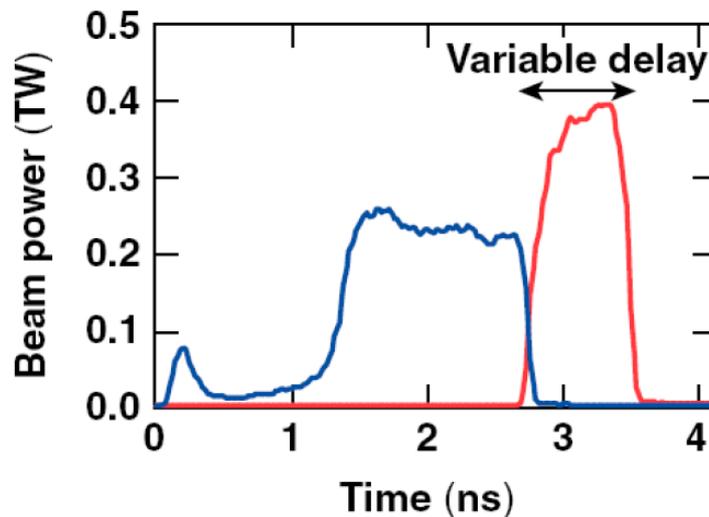
- Density scale length  $\sim 200$   $\mu$ m
- The delay and intensity of the tightly focused beams were varied
- Laser backscattering and hot-electron generation were studied

# Up to 35% of the shock-beam laser energy is lost due to backscatter



- No measurable signal of the 3/2 harmonic
- SRS dominates back reflection at highest intensity
- SBS reflection is relatively stable at ~10%

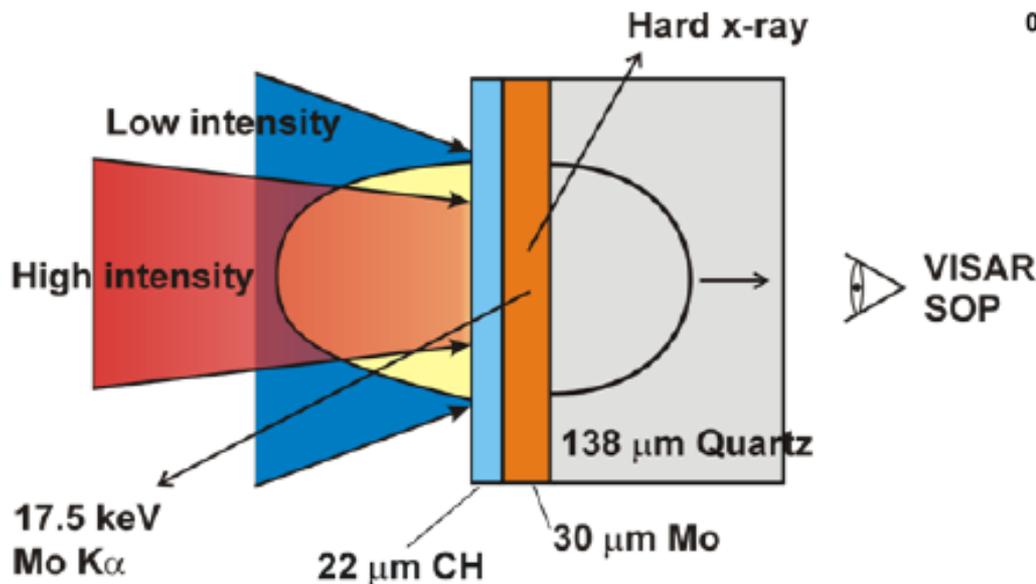
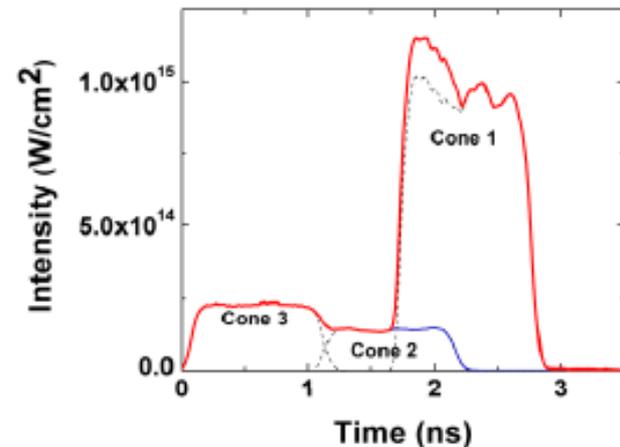
# Up to 16% of the shock-beam energy is converted into hot electrons of 45-keV temperature



# A laser-plasma interaction experiment was performed in planar geometry with overlapping beams



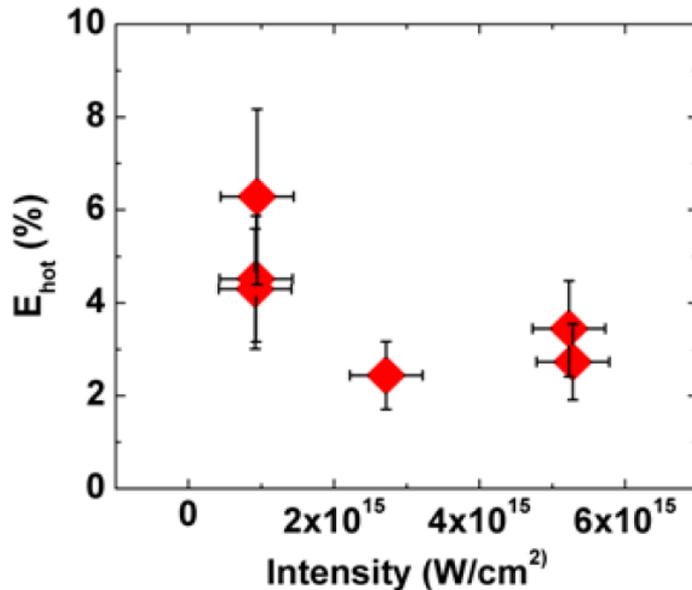
- Pre-plasma:  $\sim 2 \times 10^{14} \text{ W/cm}^2$
- Shock:  $\sim 1 - 6 \times 10^{15} \text{ W/cm}^2$
- Density scale length:  $\sim 500 \mu\text{m}$



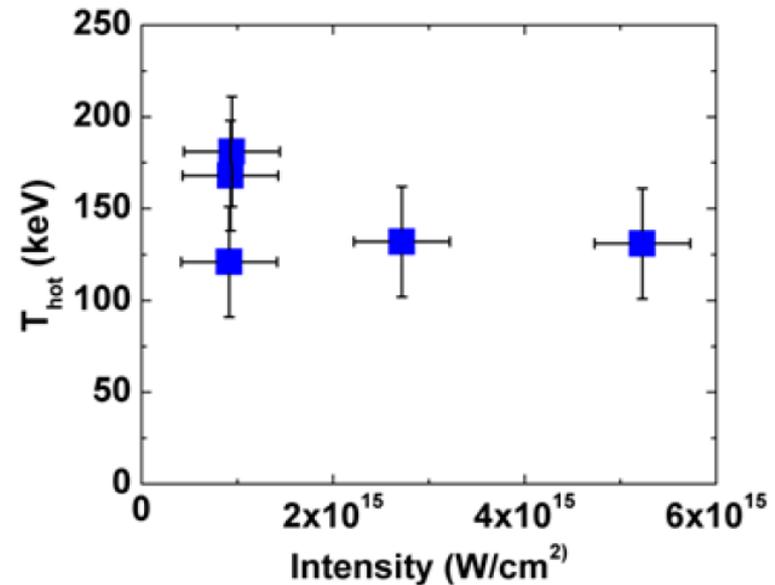
# Up to 6% of the high intensity laser energy is converted into hot electrons



Hot electron energy fraction versus intensity

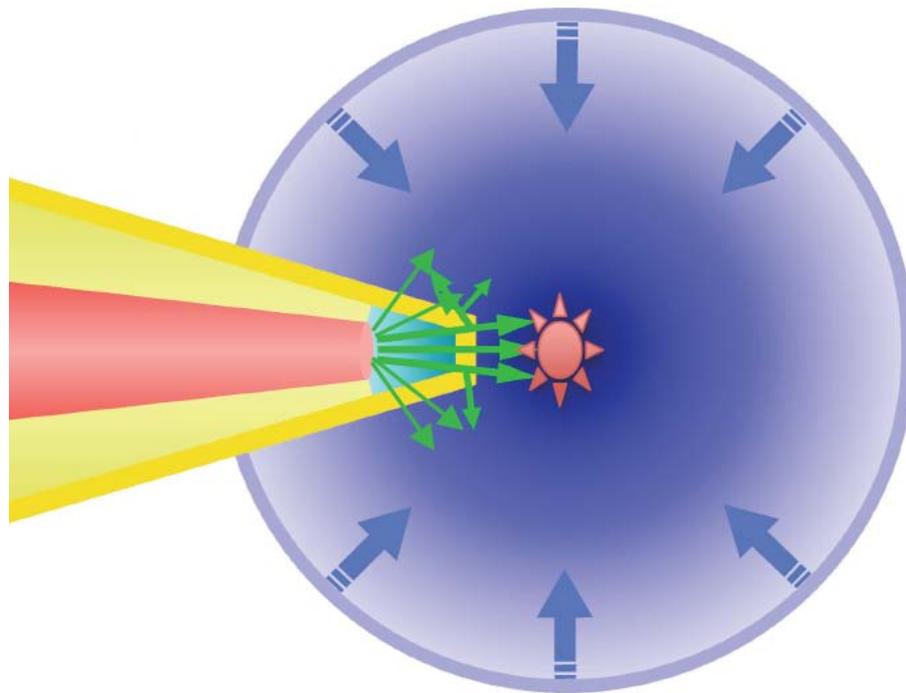


Hot electron temperature versus intensity



- The measured hot electron temperature is a factor  $\sim 3$  higher compared to spherical target experiment

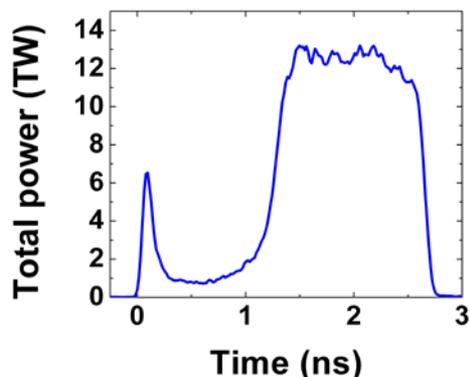
# Fast Ignition



M. Tabak et al., PoP 1, 626 (1994)

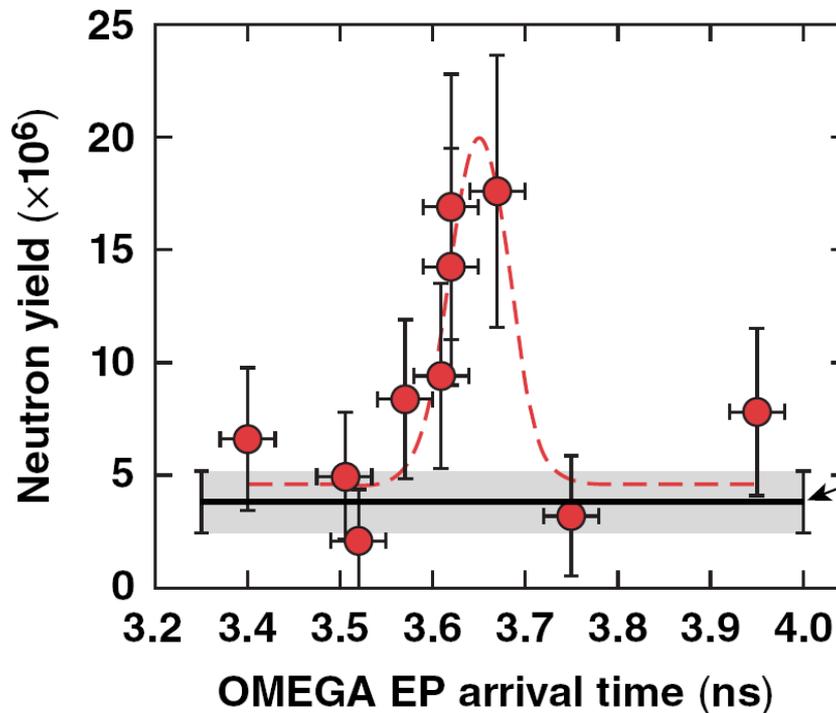
S Atzeni et al., PPCF 51, 015016 (2009)

# Fast electron heating is observed in fast ignition integrated experiments on OMEGA



Implosion

Energy	~18 kJ (54 beams)
Wavelength	351 nm
Pulse shape	Low-adiabat, $\alpha \approx 1.5$
Pulse duration	~3 ns
Implosion velocity	~ $2 \times 10^7$ cm/s



Without OMEGA EP

Heating beam

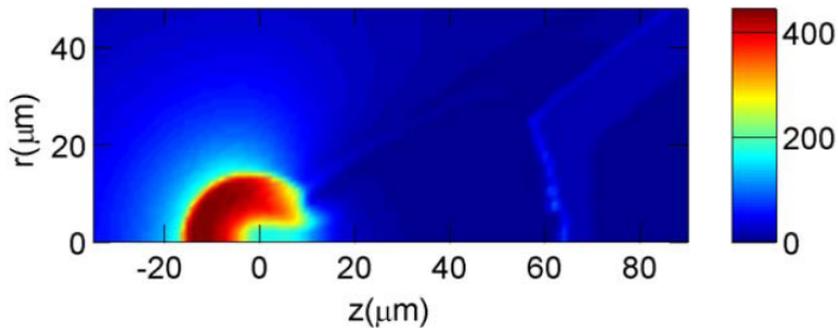
Energy	~1.0 kJ
Wavelength	1053 nm
Pulse duration	~10 ps
Intensity	~ $1 \times 10^{19}$ W/cm <sup>2</sup>

# Low-energy electrons do not heat the core in integrated DRACO-LSP simulation

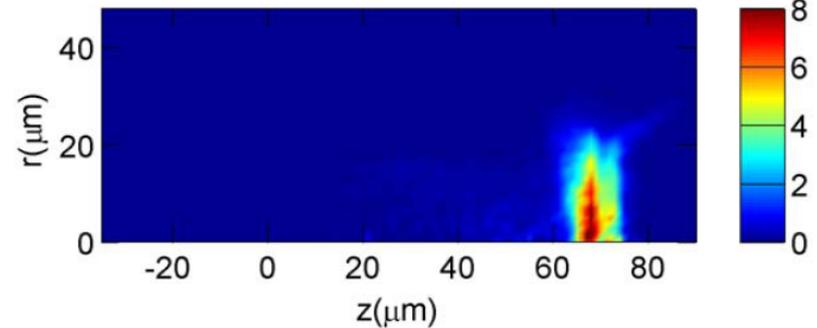


- Simulation for 10ps, 1kJ,  $R_{80}=27\mu\text{m}$ , 20% EP energy converted into fast electrons. Injection before peak  $\rho R$
- $n_{hot}$  and  $B$  are shown at the peak of the laser pulse

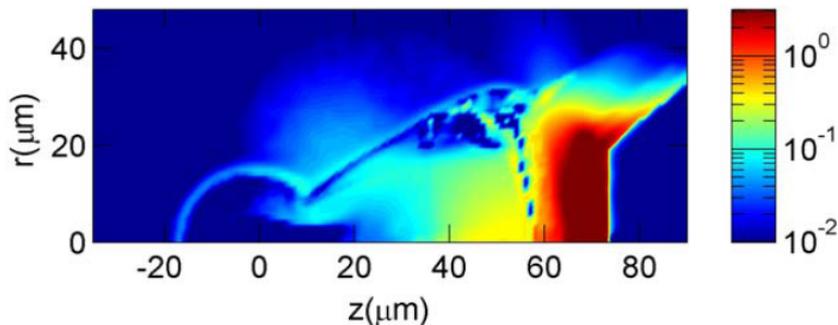
Plasma density ( $\text{g}/\text{cm}^3$ )



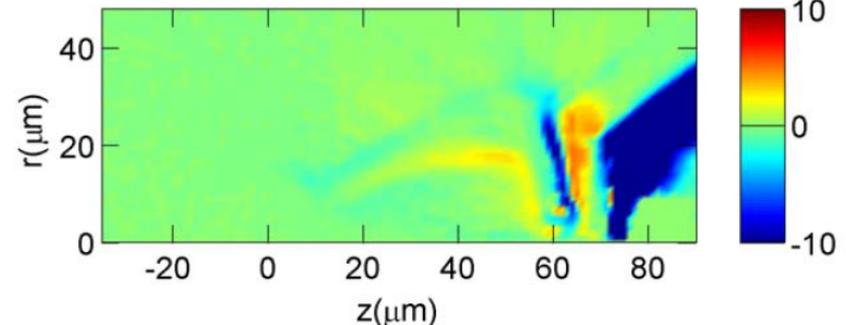
Electron beam density ( $\text{cm}^{-3} \times 10^{21}$ )



Max. plasma temperature increase (keV)



Azimuthal magnetic field (MG)



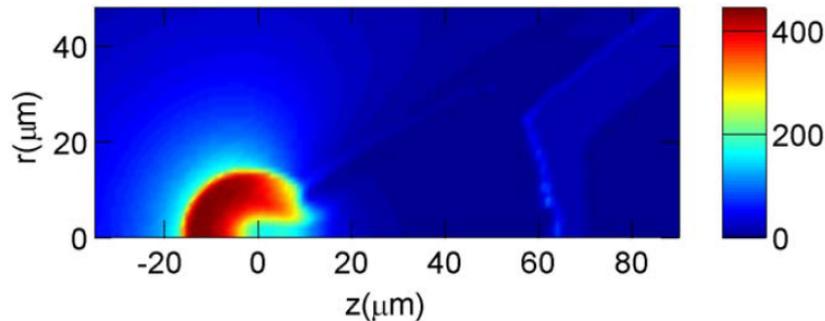
# The simulations predict an improved fast electron coupling at higher laser intensity



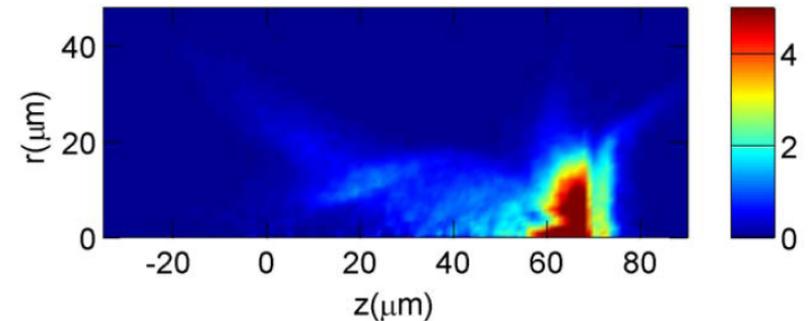
Simulation for 10ps, 2.6kJ,  $R_{80}=15\mu\text{m}$ . Injection before peak  $\rho\text{R}$

- CE ( $>100\text{ g/cm}^3$ ) improves from 0.6% to 2.4%
- CE ( $>10\text{ g/cm}^3$ ) slightly improves from 5% to 6%

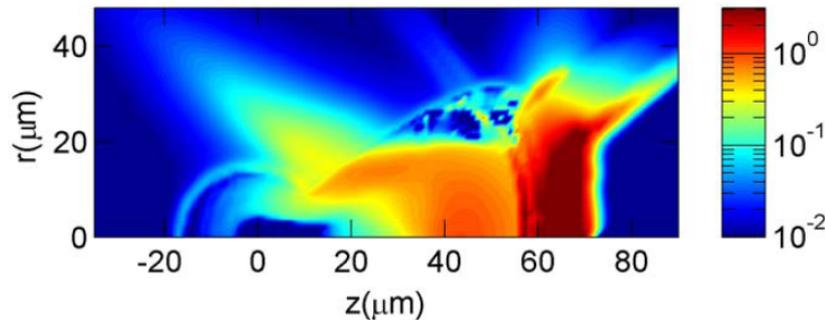
Plasma density ( $\text{g/cm}^3$ )



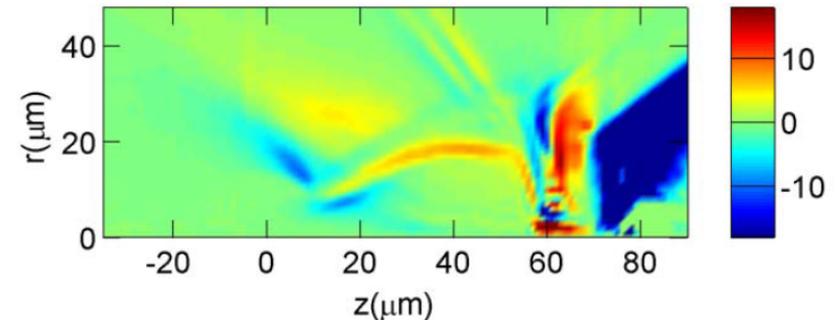
Electron beam density ( $\text{cm}^{-3}\times 10^{21}$ )



Max. plasma temperature increase (keV)



Azimuthal magnetic field (MG)

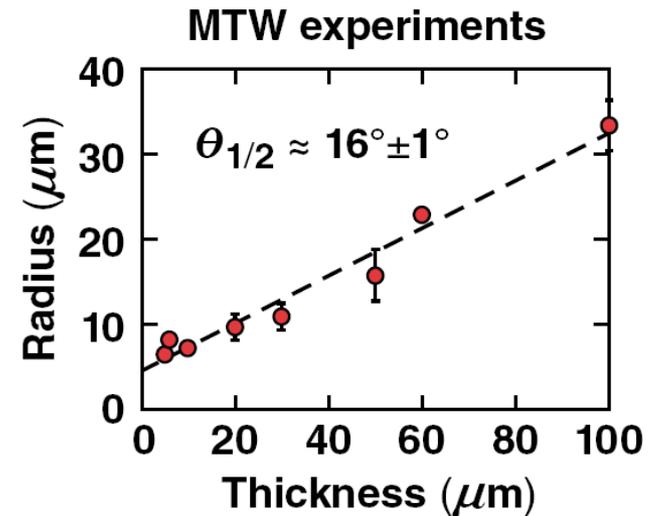
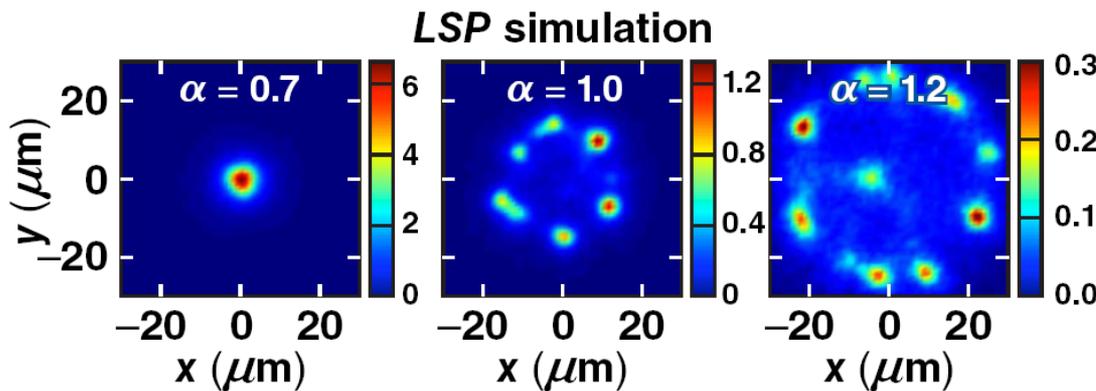
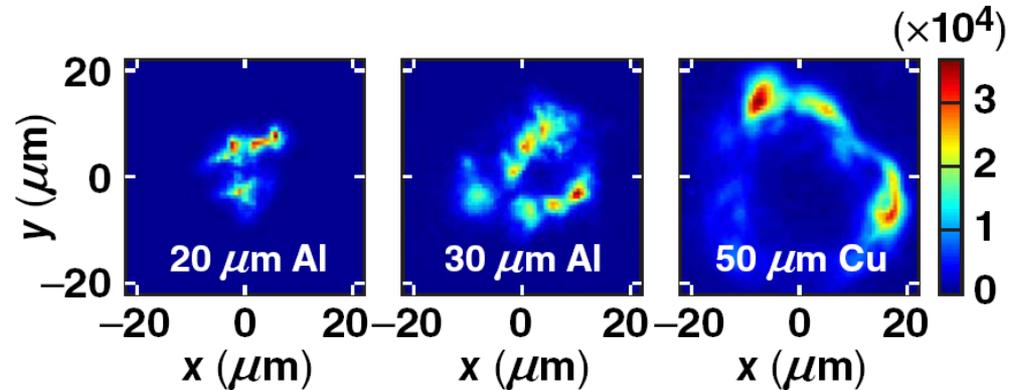
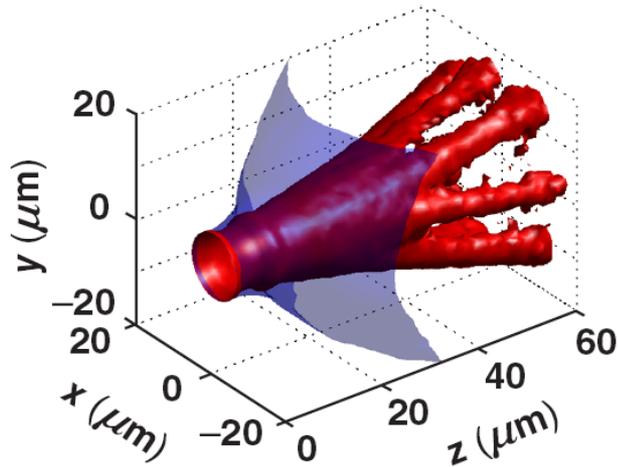


# About 60 researchers participated to the 9th FSC meeting at LLNL—a special two-day topical meeting was devoted to assess electron divergence in fast ignition



The meeting produced a final report available on the FSC website.

# Experiments on MTW and LSP simulations study fast-electron divergence and magnetic collimation



# Various type of targets were used to compare fast electrons transport in partially and fully driven, un-driven foams and cold CH

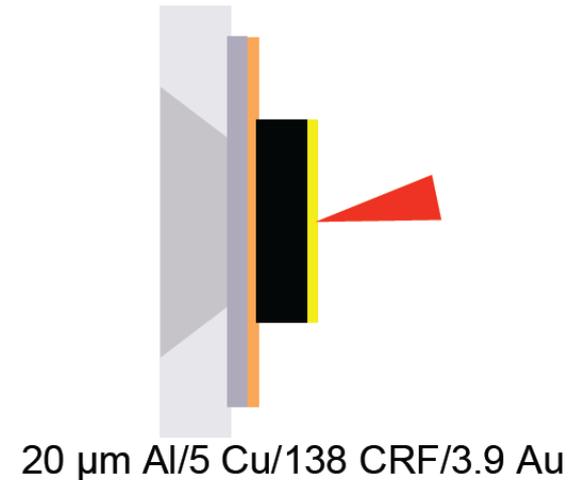
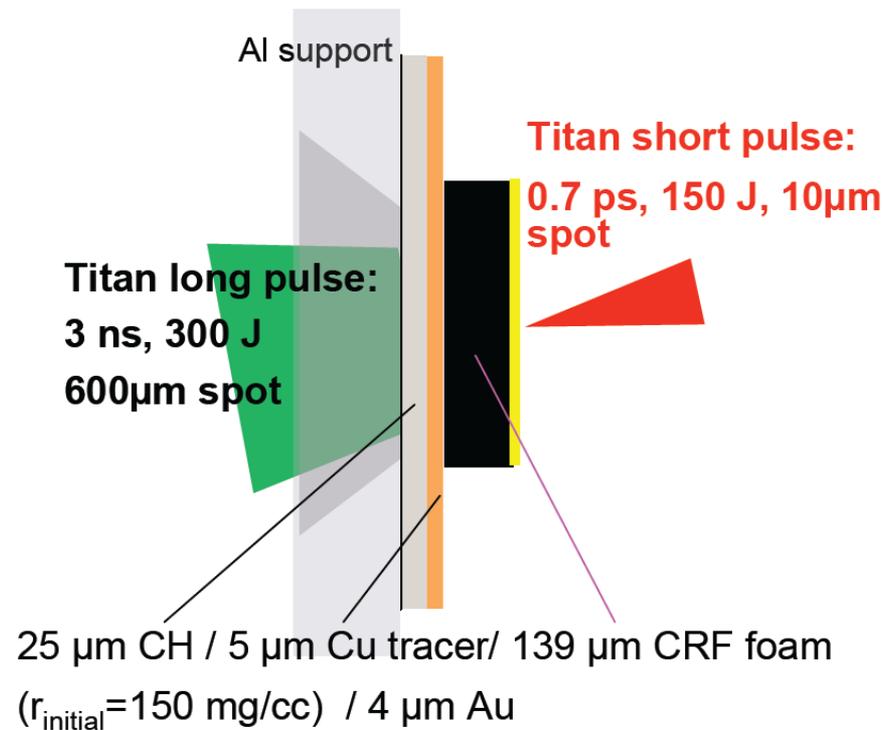


GENERAL ATOMICS

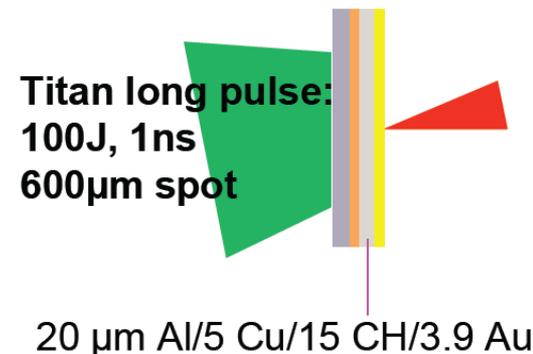


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## Foam package targets for hot transport (with various SP and LP timing delays)



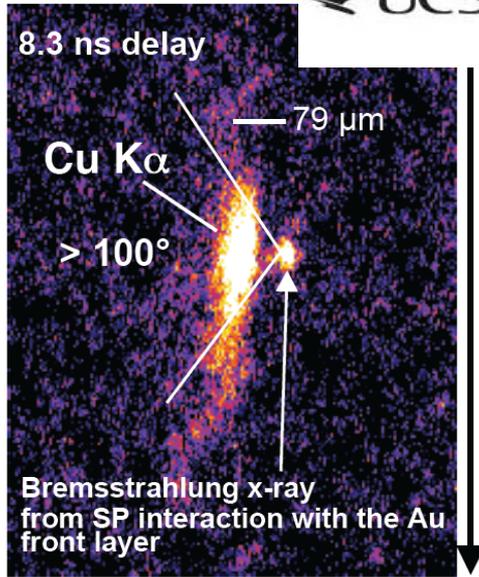
## CH insulator as transport layer



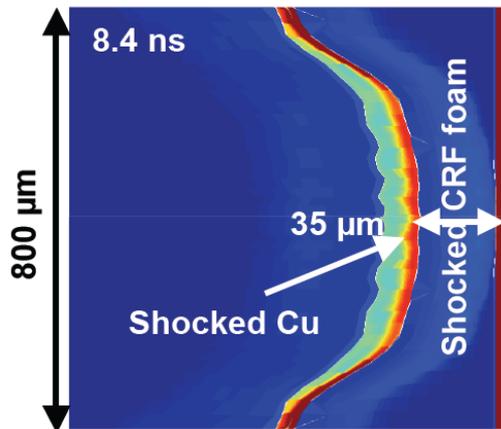
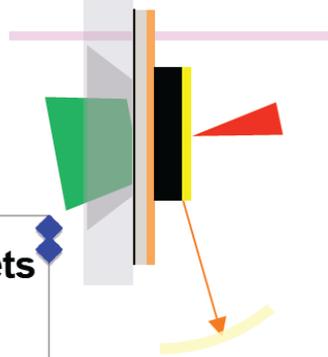
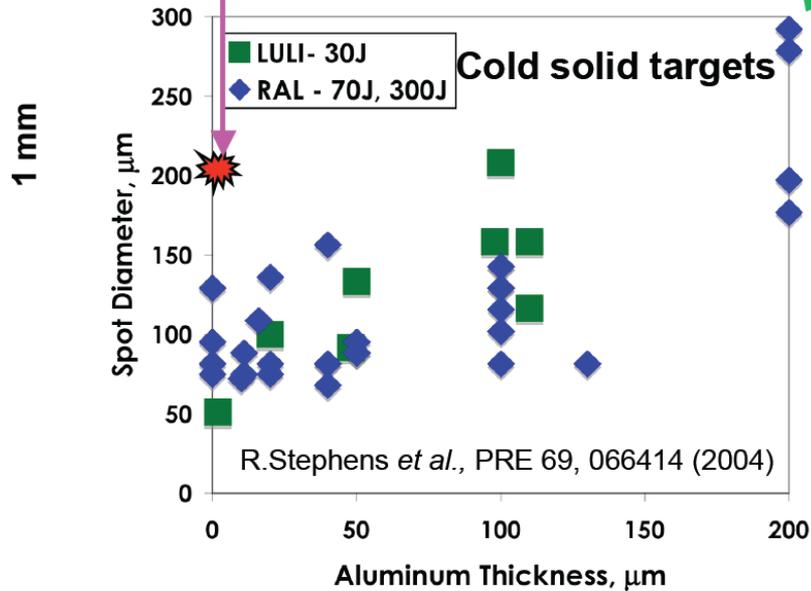
# Large extended Cu K $\alpha$ spot was consistently observed in WDM targets suggesting a large angular spread of fast electrons



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## New results from WDM experiments



- 2X larger K $\alpha$  spot in WDM case compared to the cold solid targets results

- Such large extended K $\alpha$  spot was neither observed in un-driven and partially driven (at 3 ns delay case) foam targets, nor in CH insulator targets

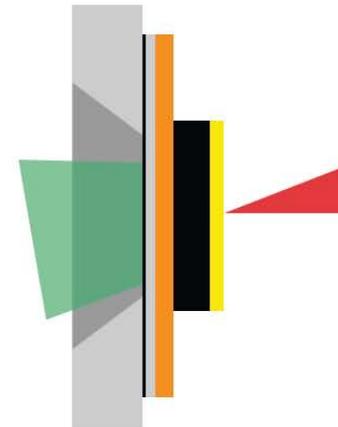
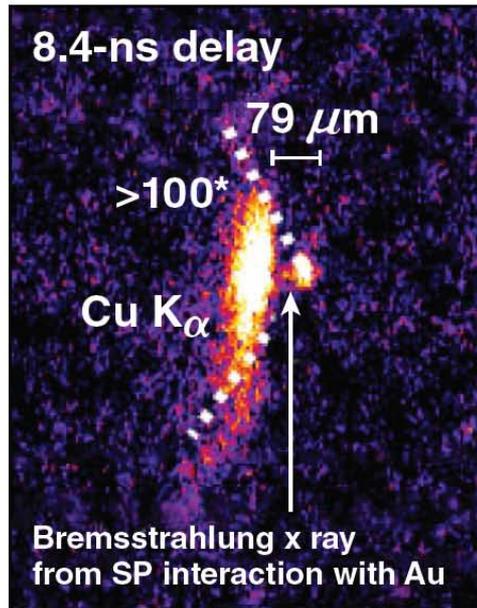
# The NIF can explore advanced ignition options



- **With day-one hardware, the NIF can explore high-gain shock ignition**
  - **Polar Shock Ignition (uses half the NIF beams to drive the implosion and the other half to drive the ignitor shock)**
- **Fast Ignition requires major hardware upgrades: 100kJ-class multi-PW laser [also talk by P. Patel at this meeting]**
- **Polar Direct Drive requires minor upgrades: multi-FM or 2D-SSD (talk by J. Soures at this meeting)**

**Back-up slides**

# Fast-electron transport in WDM is investigated in a set of experiments on Titan using foam targets



A large angular spread ( $>90$ ) has been inferred from the size of the Cu  $K_{\alpha}$  emission spot.

**About 3% of the electron-beam energy is deposited in the core region with  $\rho > 100 \text{ g/cm}^3$**



**Energy deposition**

	<b>Fraction of e-beam energy</b>	<b>Fraction of laser energy</b>
<b>Deposition in gold</b>	<b>52%</b>	<b>10%</b>
<b>Deposition in plastic with <math>\rho &gt; 10 \text{ g/cm}^3</math></b>	<b>25%</b>	<b>5%</b>
<b>Deposition in plastic with <math>\rho &gt; 100 \text{ g/cm}^3</math></b>	<b>3%</b>	<b>0.6 %</b>

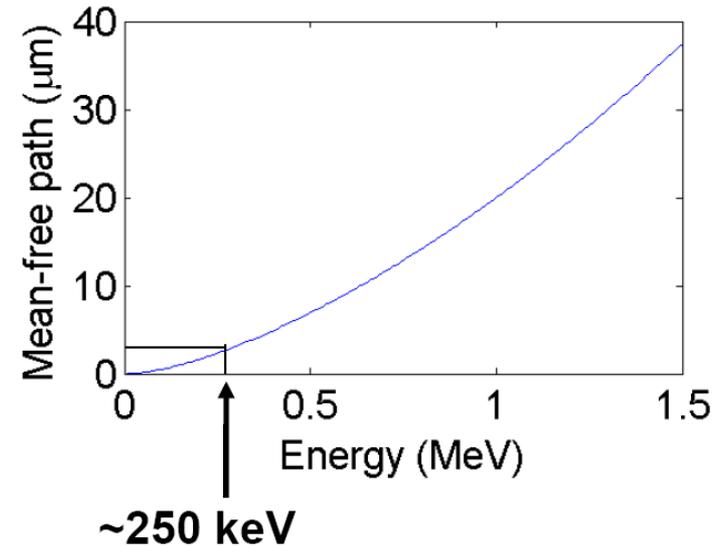
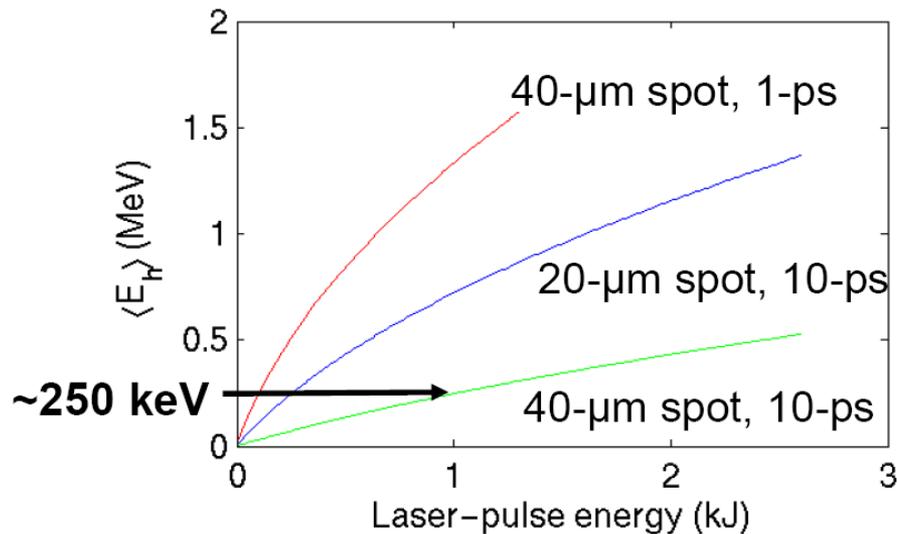
**Neutron yield increase**

<b>Neutron yield without hot electrons</b>	<b><math>6.6 \times 10^8</math></b>
<b>Neutron yield with hot electrons</b>	<b><math>7.4 \times 10^8</math></b>
<b>Neutron yield increase</b>	<b><math>8 \times 10^7</math></b>
<b>Neutron yield increase in the region with <math>\rho &gt; 100 \text{ g/cm}^3</math></b>	<b><math>1.6 \times 10^7</math></b>

# The hot-electron energy can be too low for a good penetration through the Au cone tip

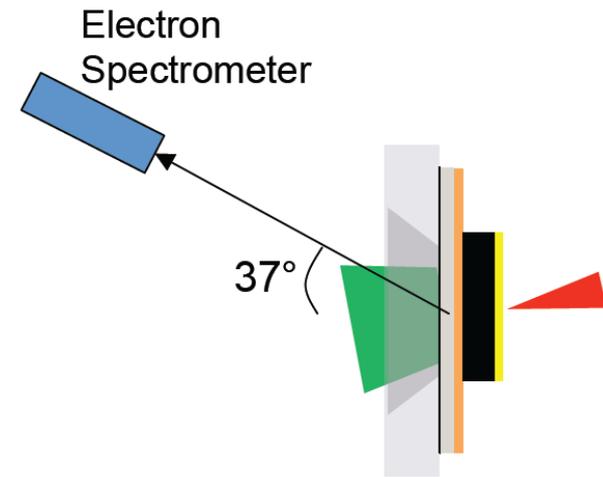
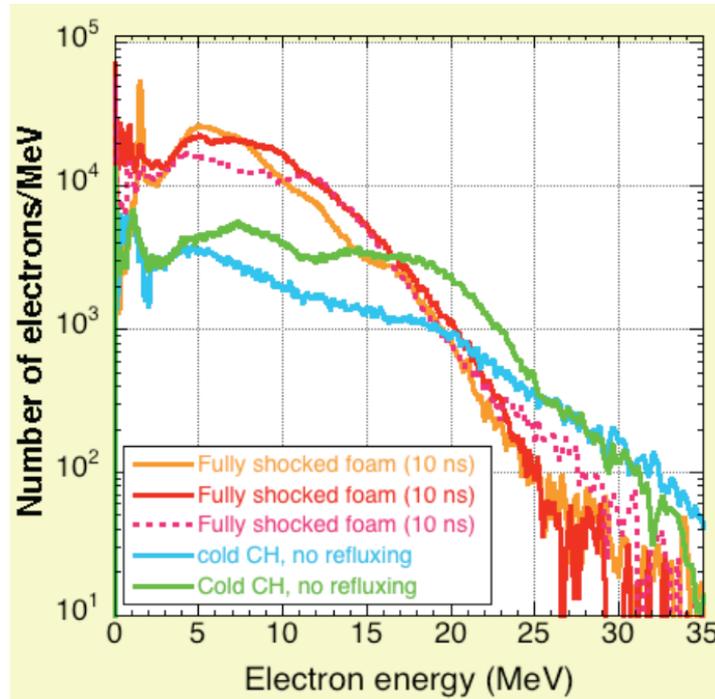


Mean hot-electron energy assuming ponderomotive scaling  
(averaged within FWHM of the spatial and temporal distribution for a Gaussian pulse)



- Mean-free path of 250 keV electrons is a few  $\mu\text{m}$  and is smaller than the cone wall thickness
- Higher laser intensities are required

## 5X more electrons were emitted sideways in WDM compared to CH insulator case



- Significant increase in the number of escaped electrons from the side is consistent with large angular spread of electrons in WDM